

氫硼核融合實作

Practical course on proton-boron nuclear fusion



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Institute of Space and Plasma Sciences, National Cheng Kung University

2025 summer break

7/14(Mon.) – 7/18(Fri.) 14:00-17:40

Lecture 2

<https://capst.ncku.edu.tw/PGS/index.php/teaching/>

Course Outline



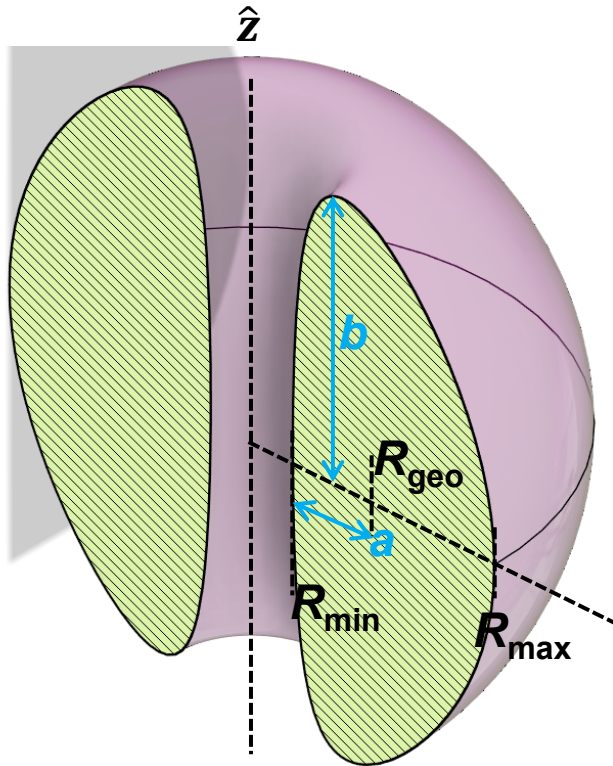
- **Formosa Integrated Research Spherical Tokamak (FIRST), the first Tokamak in Taiwan**
- **Physics used in Alpha-E**
 - **Microwave**
 - **ECRH**
 - **Spectrum of fusion products**
 - **Stopping power of the filter**
 - **Data analysis using Pulse Shaping**
- **Final report**

Course Outline



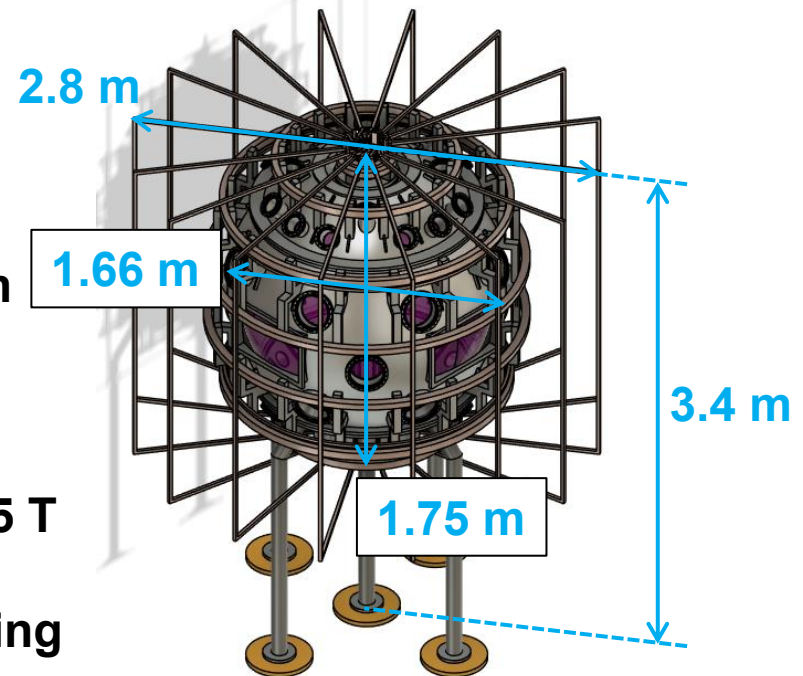
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Development of Formosa Integrated Research Spherical Tokamak (FIRST), the first Tokamak in Taiwan

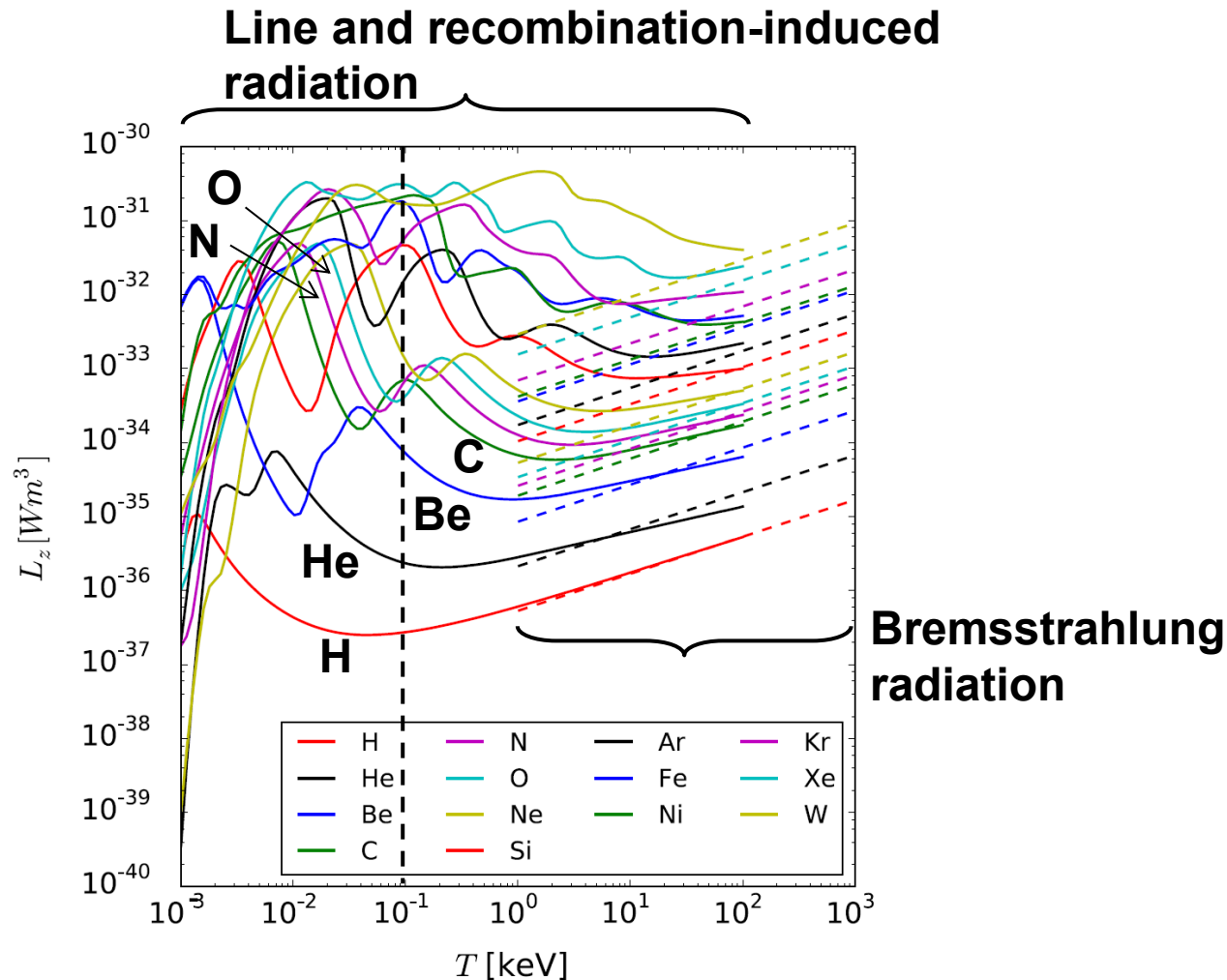


- $R_{\text{geo}} = 45 \text{ cm}$
- $a = 32 \text{ cm}$
- $b = 76.8 \text{ cm}$
- $T \sim 100 \text{ eV}$
- $B_T \sim 0.1 - 0.5 \text{ T}$
- $I_p \sim 100 \text{ kA}$
- Ohmic heating
- Gas: H_2
- Duration: 100 ms

• **First plasma: 2026.**



Temperature of 100 eV is the threshold of radiation barrier by impurities



Teams



- 馬維揚團隊 @ 國家原子能科技研究院
 - Site
- 張博宇團隊 @ 成功大學 太空與電漿科學研究所
 - System design and development/diagnostics
- 向克強團隊 @ 成功大學 前瞻電漿中心
 - Theoretical design
- 河森榮一郎團隊 @ 成功大學 太空與電漿科學研究所
 - Diagnostics
- 柳克強團隊 @ 清華大學 工程與系統科學系
 - Diagnostics
- 蔡宗哲團隊 @ 國家高速網路與計算中心
 - Simulation
- 張存續團隊 @ 清華大學 物理系
 - RF startup

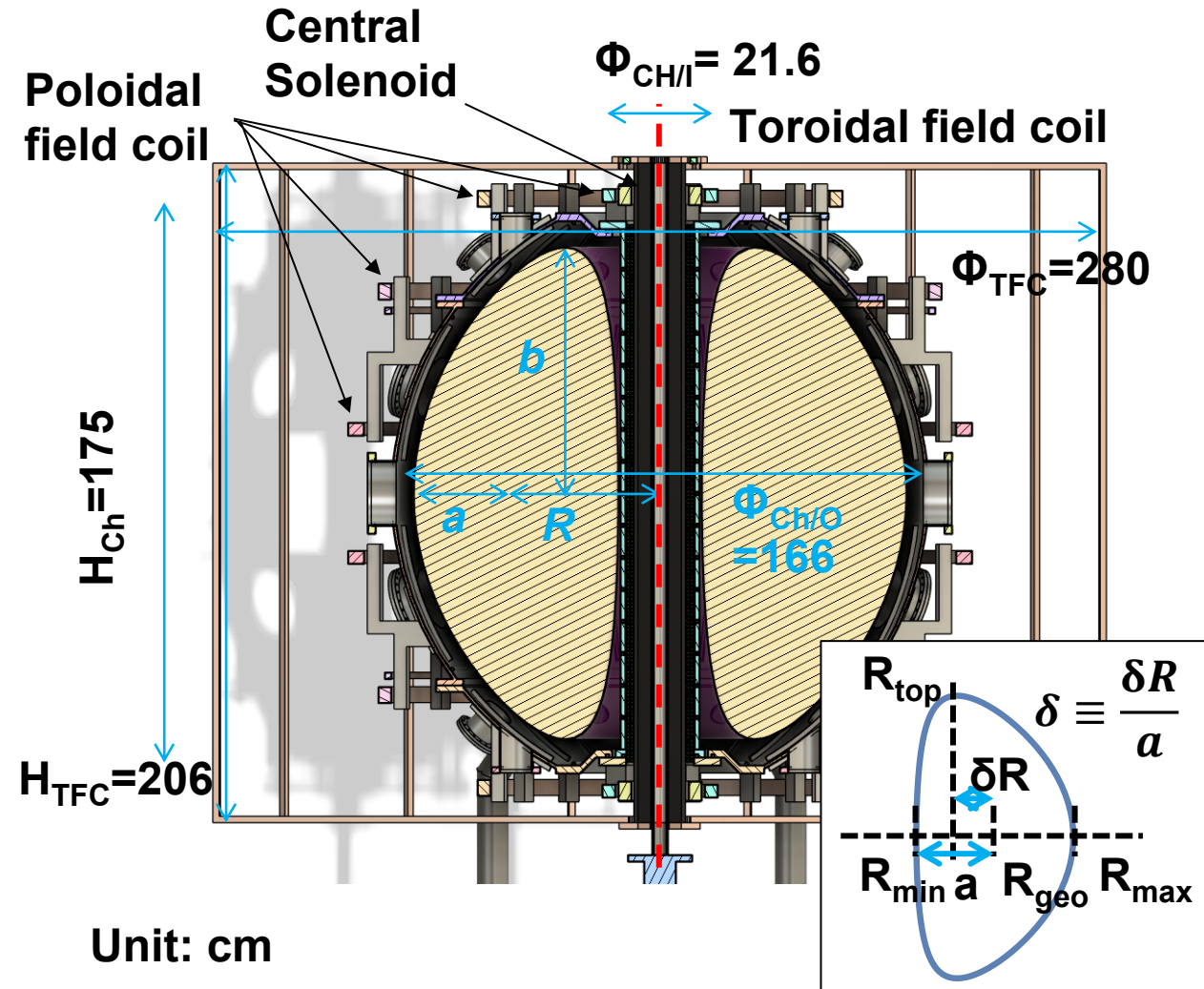
• We welcome anyone interested in fusion research to join us!

The system will be built in National Atomic Research Institute (NARI, 國家原子能科技研究院) at 龍潭



https://www.sipa.gov.tw/home.jsp?mserno=201001210037&serno=201001210041&menudata=ChineseMenu&contlink=content/introduction_4_1.jsp&serno3=201002010023
<https://www.nari.org.tw/newsdetail/activity/353.html>

Formosa Integrated Research Spherical Tokamak (FIRST) aiming for the first plasma in 2026



- $R/a/b = 45/32/76.8$ cm
- Aspect ratio $= R/a = 1.5$
- $k = b/a = 2.4$
- $\delta = 0.5$
- $B_T = 0.1 \sim 0.5$ T
- $T \geq 100$ eV
- $I_p \geq 100$ kA
- Gas: H_2
- Ohmic heating
- Duration: 100 ms
- FIRST is targeted for
 - Low aspect ratio
 - High beta
 - High bootstrap current

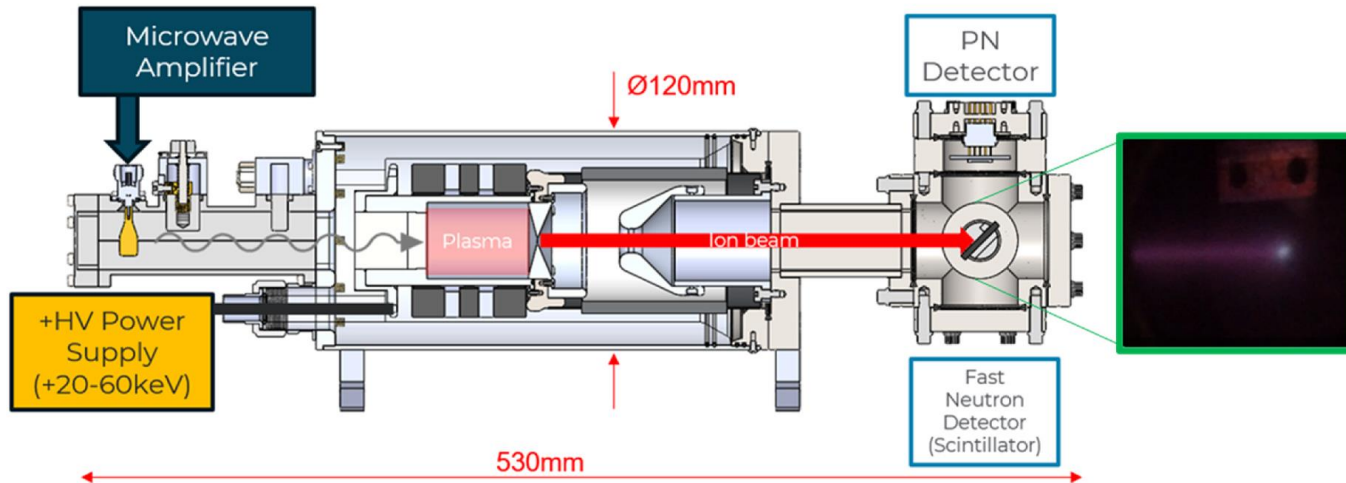
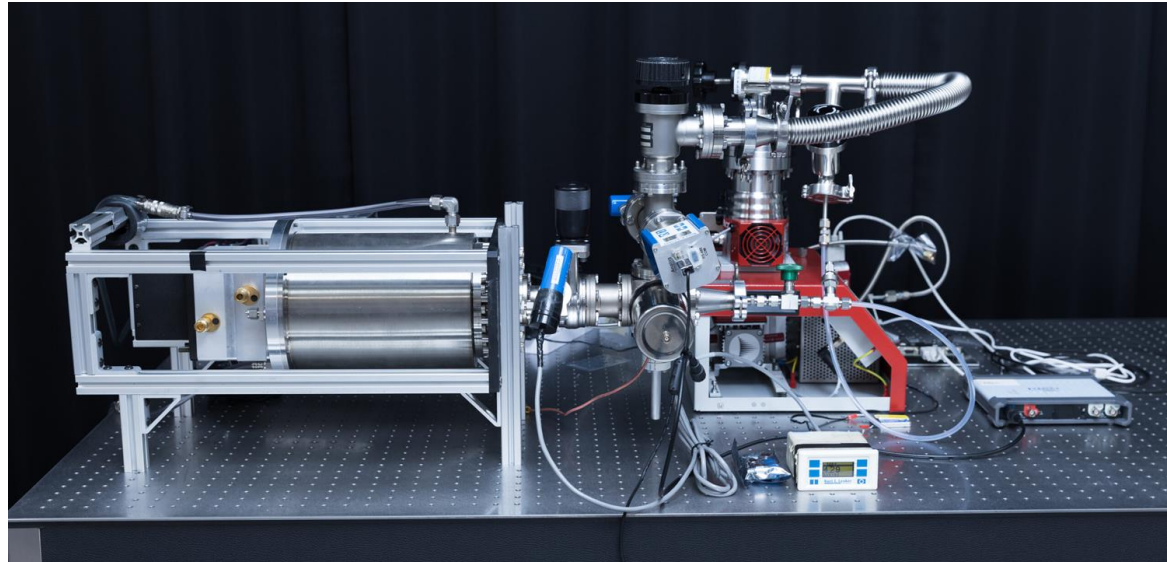
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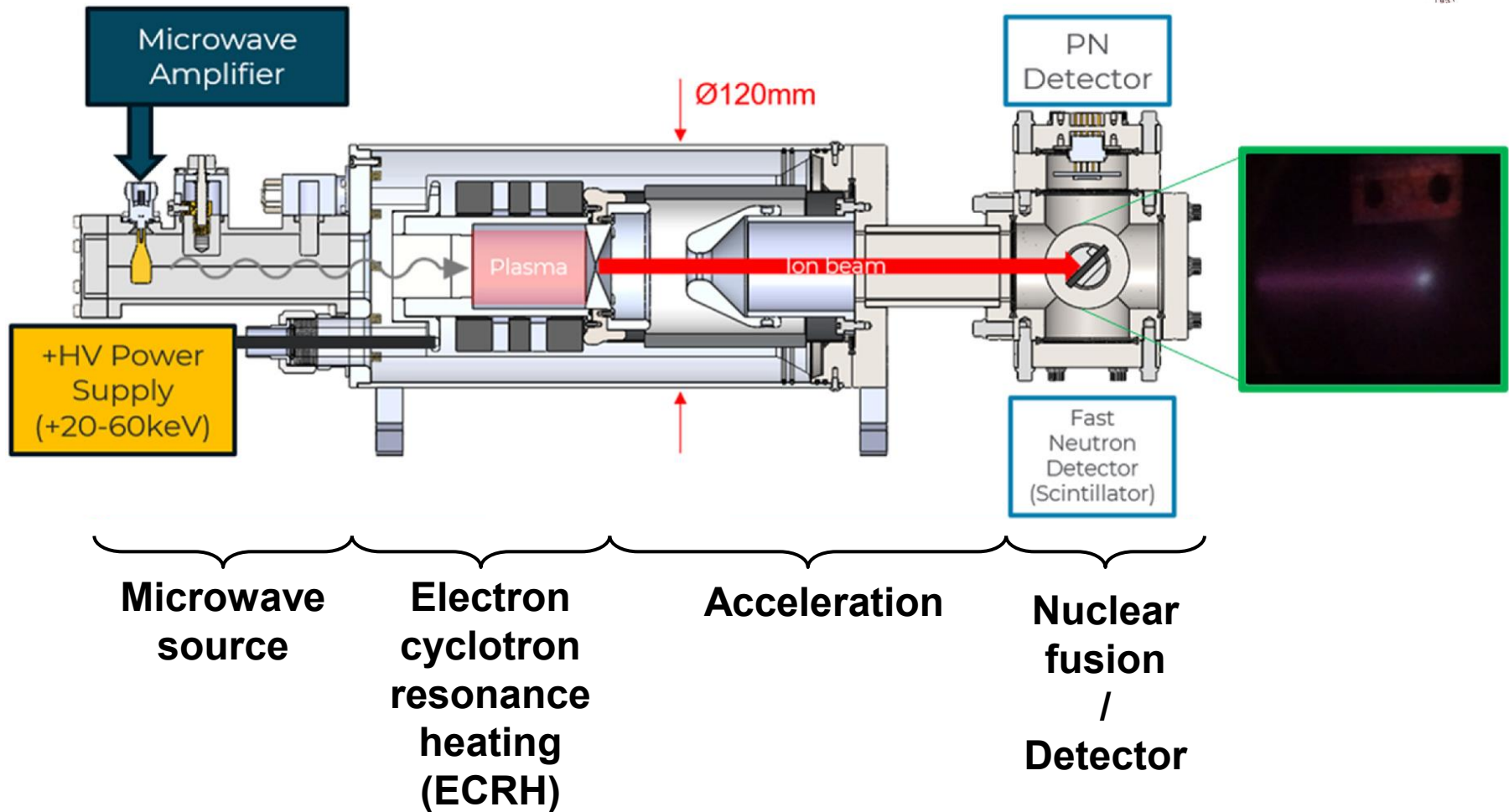
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- Final report

Alpha-E provided by Alpha Ring will be used to demonstrate the proton-boron nuclear fusion



<https://alpharing.com/alpha-e/>
Allan Xi Chen, etc., Physics Open 21, 100234(2024)

Alpha-E provided by Alpha Ring will be used to demonstrate the proton-boron nuclear fusion

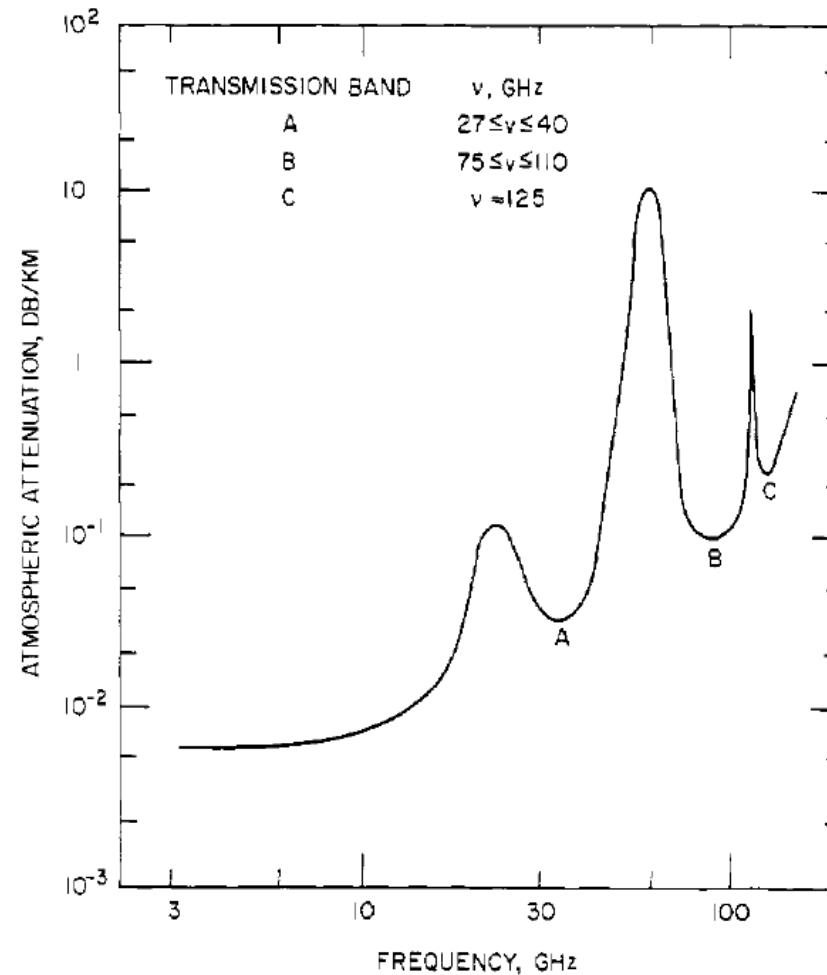


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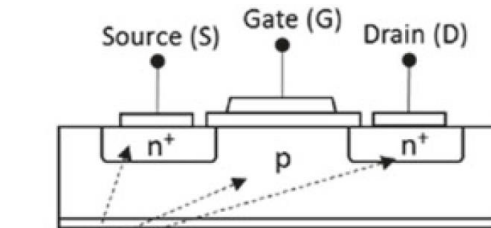
Microwave frequency is determined for those used in communications and radar purposes



Microwave / Radio-frequency (RF) wave can be generated using solid-state device



(a) Solid State Device

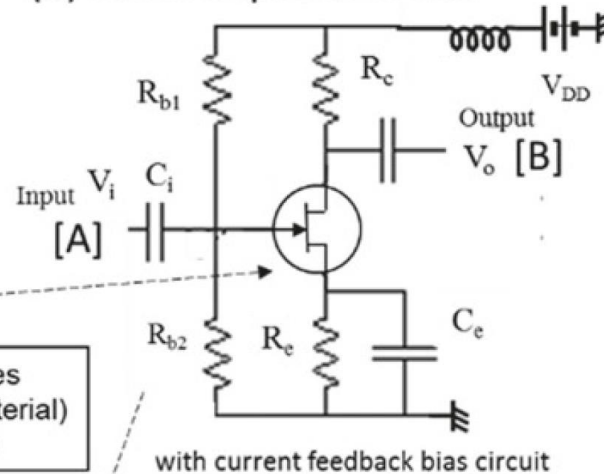


e.g. GaN
(Material)

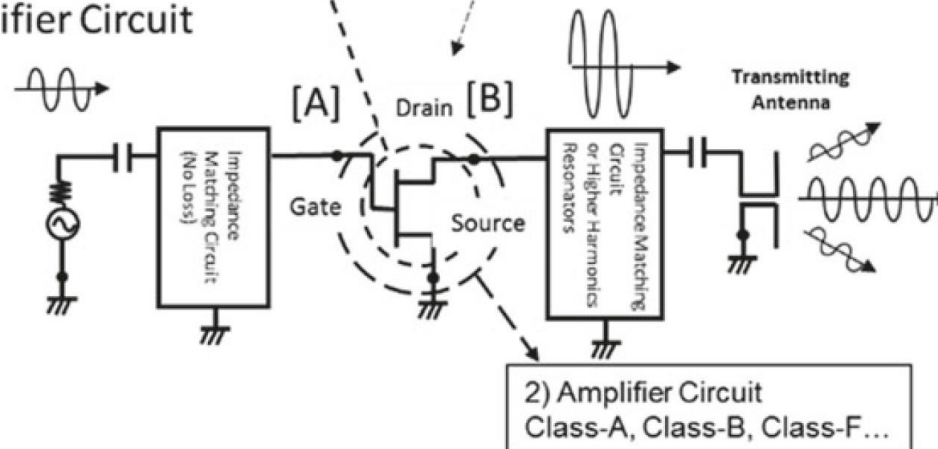
FET

1) Semiconductor Devices
Si, GaAs, SiC, GaN (Material)
FET, HBT, HEMT (Form)

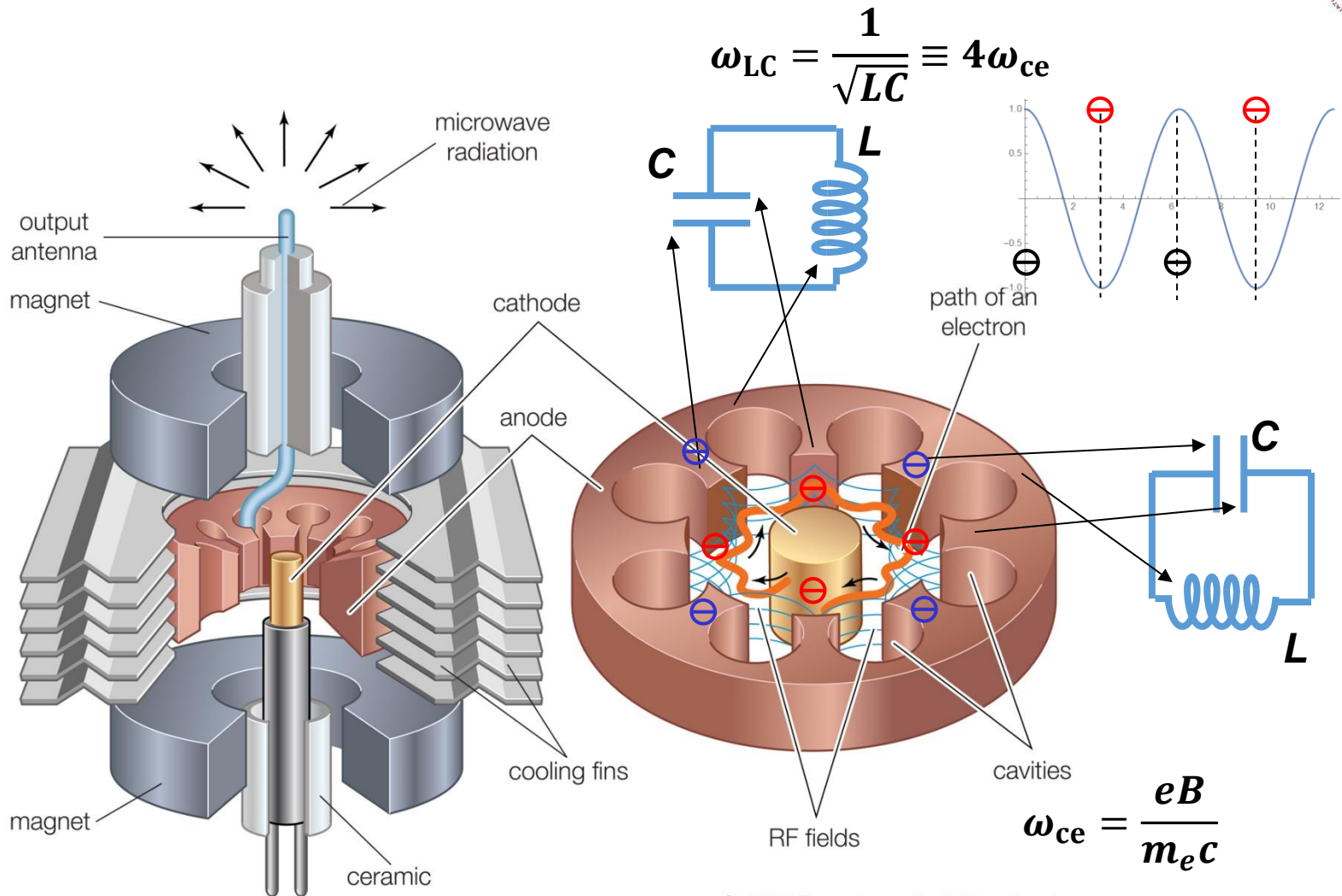
(b) Basic Amplifier Circuit



(c) Amplifier Circuit

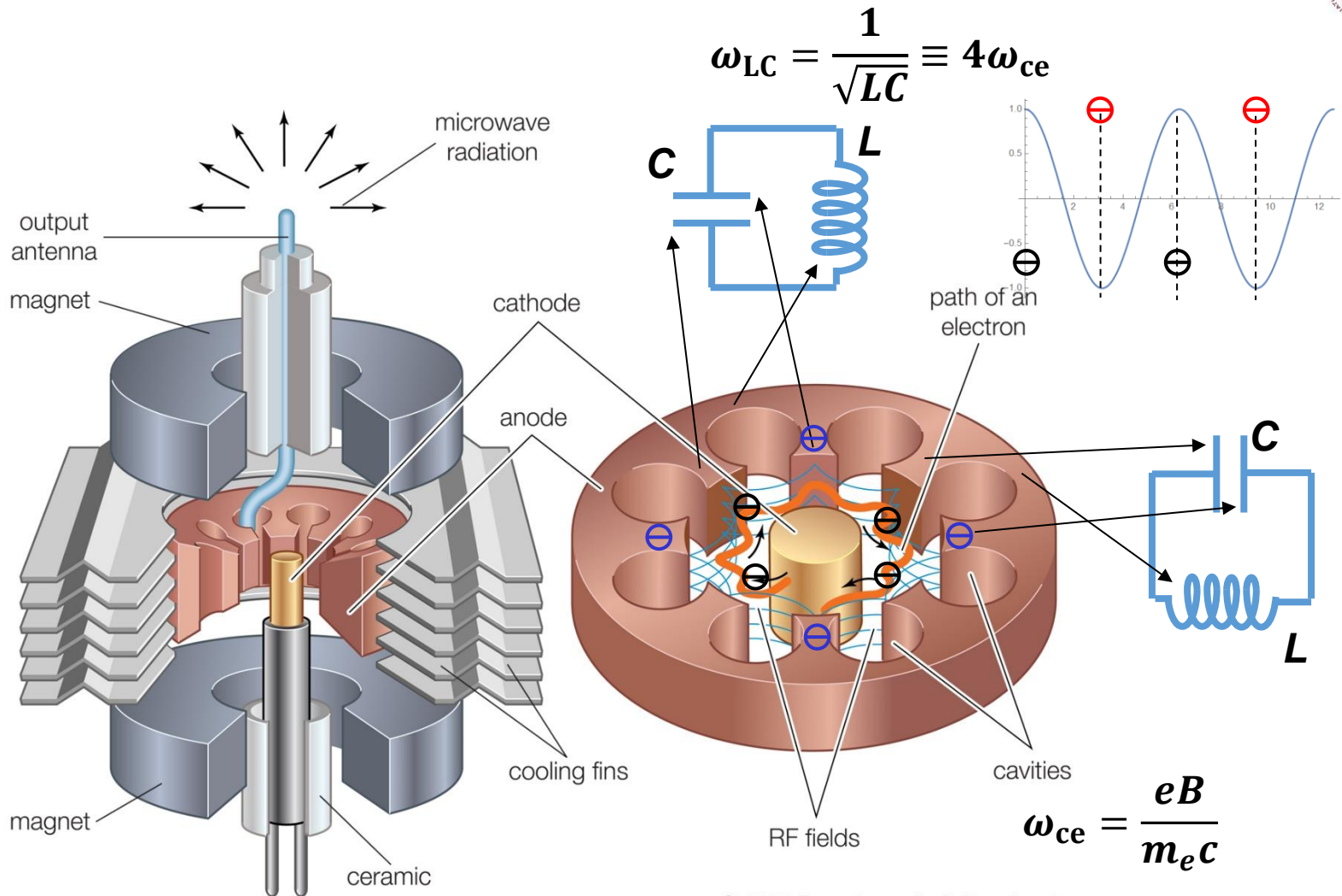


Internal of a magnetron



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Internal of a magnetron

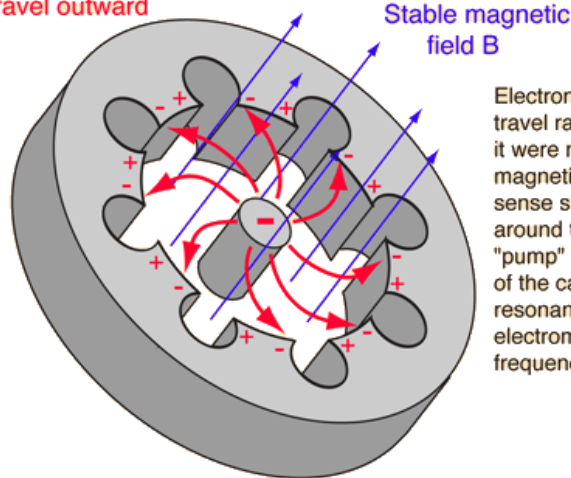


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Magnetron is a forced oscillation driven by electrons between the gap

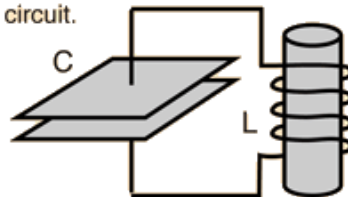


Hot cathode emits electrons which travel outward



Electrons from a hot filament would travel radially to the outside ring if it were not for the magnetic field. The magnetic force deflects them in the sense shown and they tend to sweep around the circle. In so doing, they "pump" the natural resonant frequency of the cavities. The currents around the resonant cavities cause them to radiate electromagnetic energy at that resonant frequency.

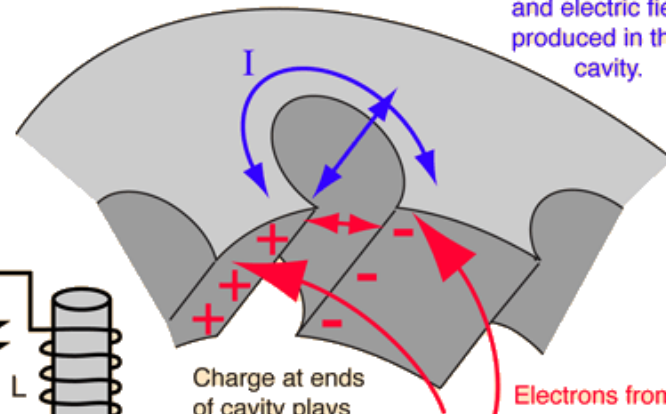
The cavity exhibits a resonance analogous to a parallel resonant circuit.



$$f_{\text{resonance}} \approx \frac{1}{2\pi} \sqrt{\frac{1}{LC}}$$

Current around the cavity plays the role of an inductor.

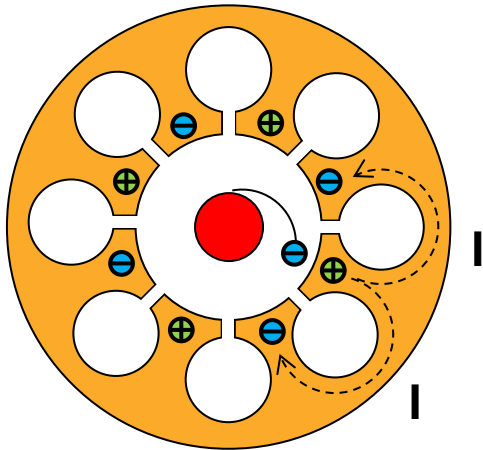
Oscillating magnetic and electric fields produced in the cavity.



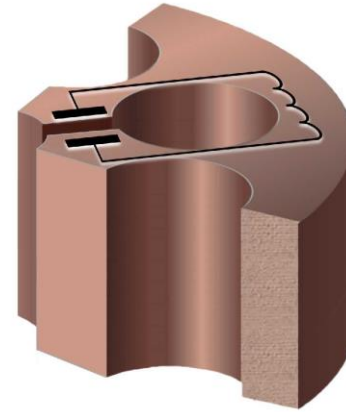
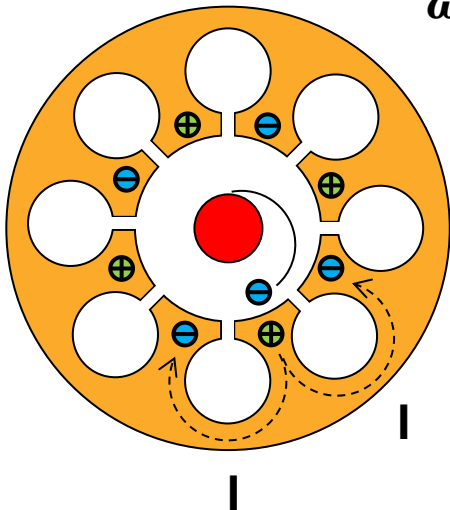
Charge at ends of cavity plays the role of a capacitor.

Electrons from the hot center cathode arriving at a negatively charged region tend to drive it back around the cavity, "pumping" the natural resonant frequency.

Strong oscillation occurs when the electron cyclotron frequency match the LC oscillation frequency



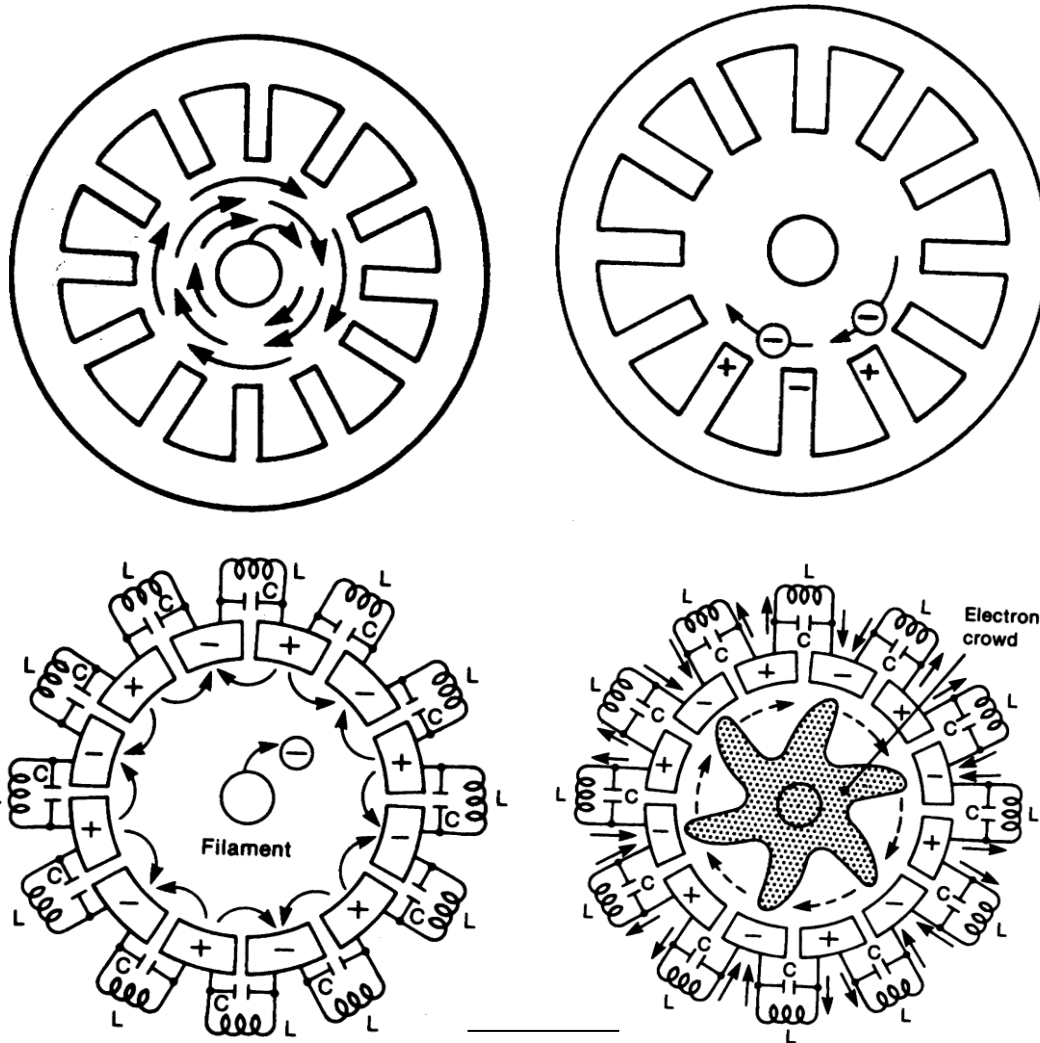
$$\omega_{\text{CE}} = \frac{eB}{mc}$$



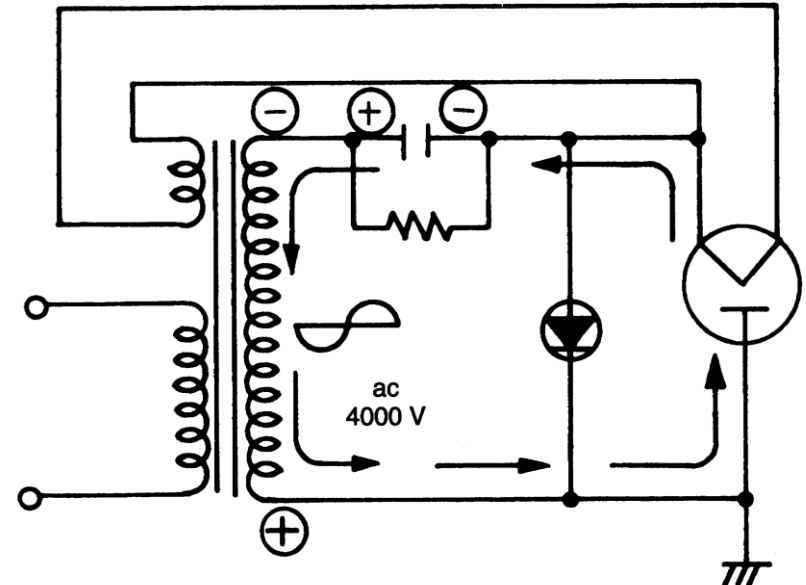
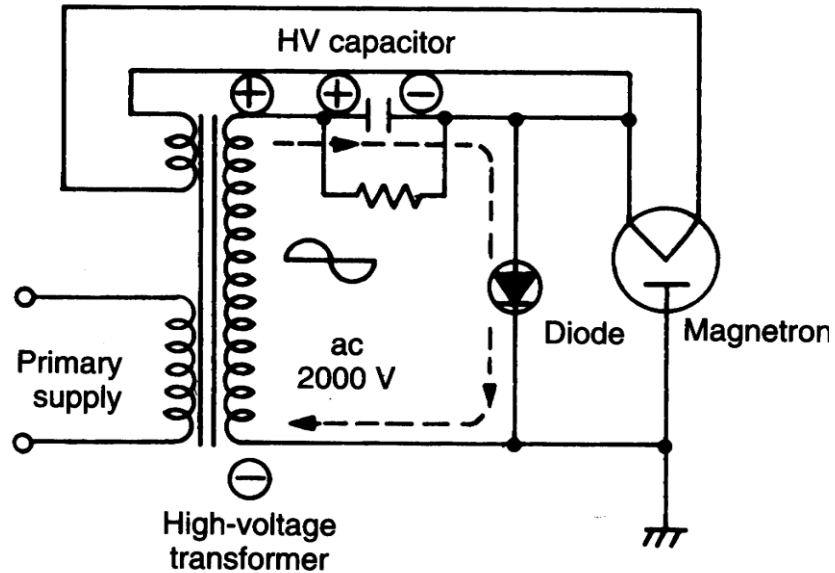
$$\omega = \frac{1}{\sqrt{LC}}$$

Resonance condition: $\omega_{\text{CE}} = \omega$

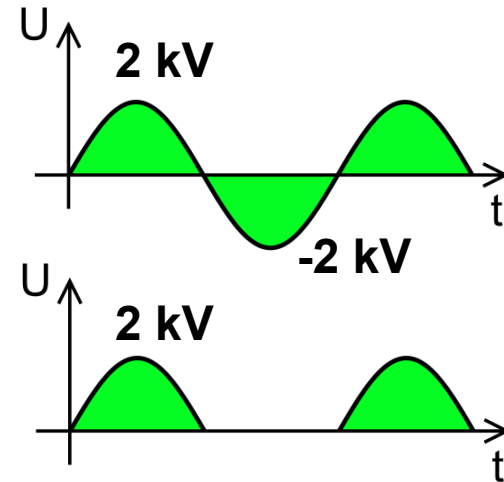
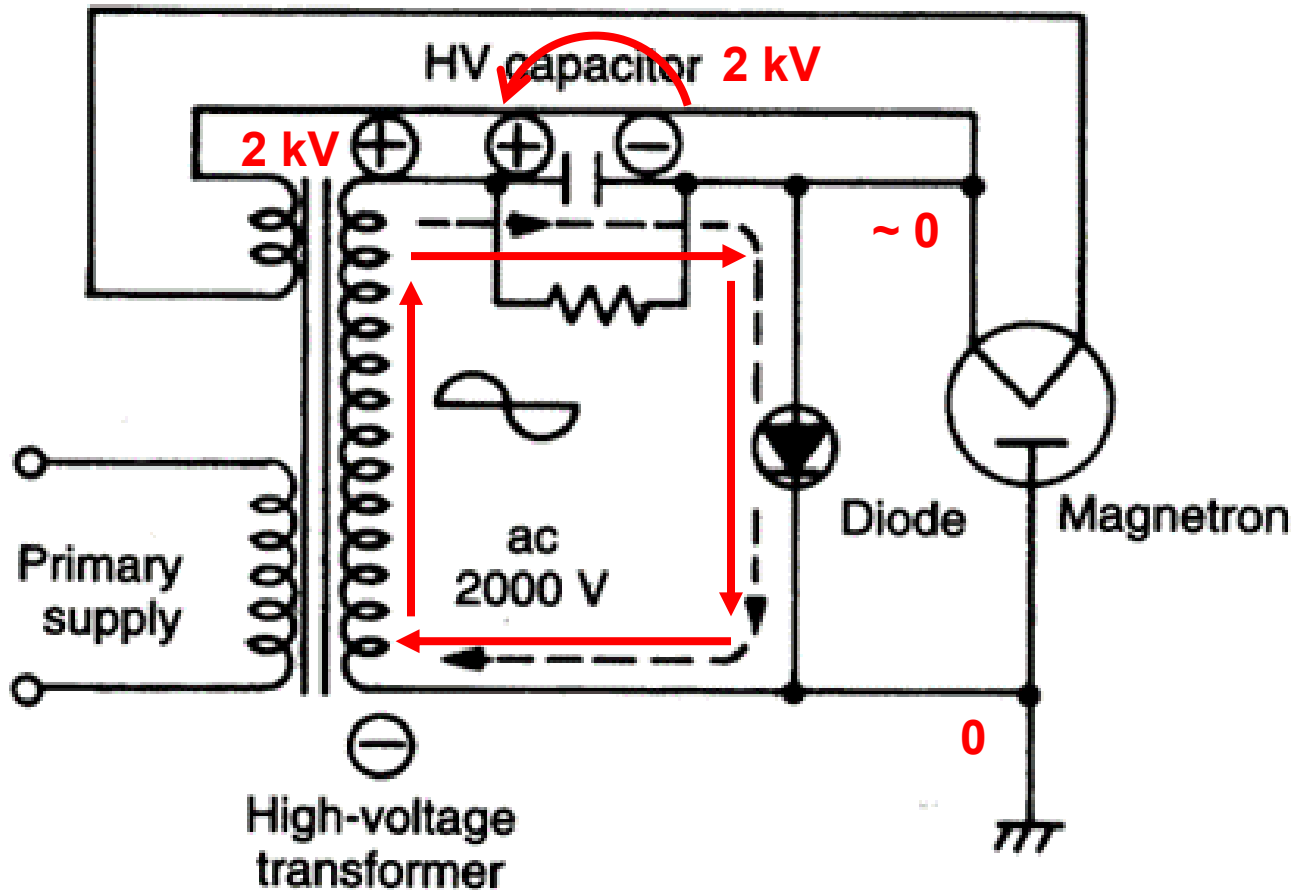
Resonance in a magnetron



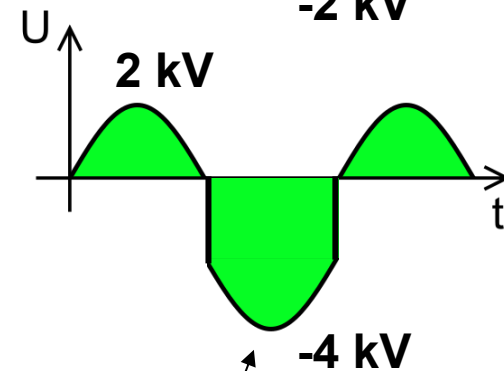
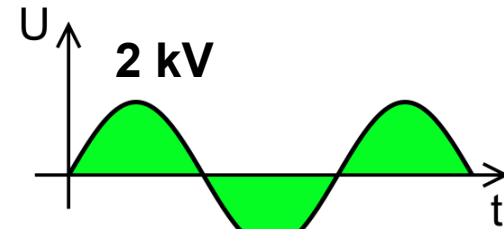
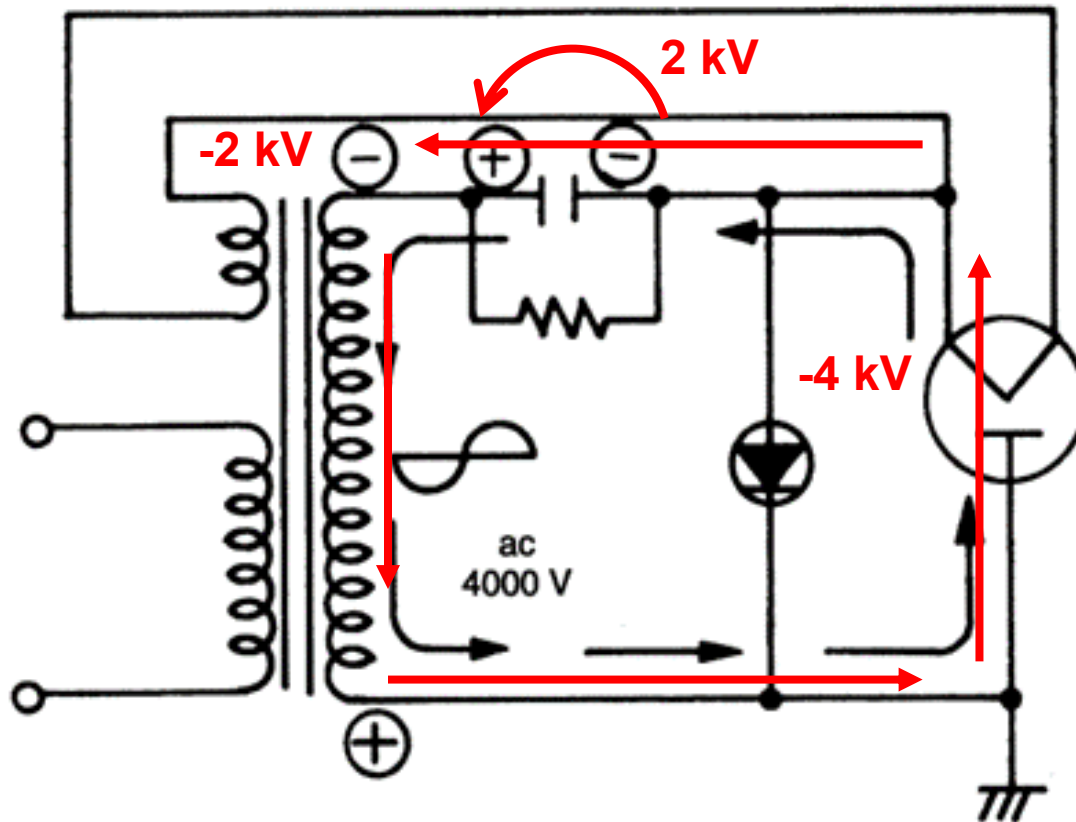
Magnetron schematic diagram



Magnetron schematic diagram

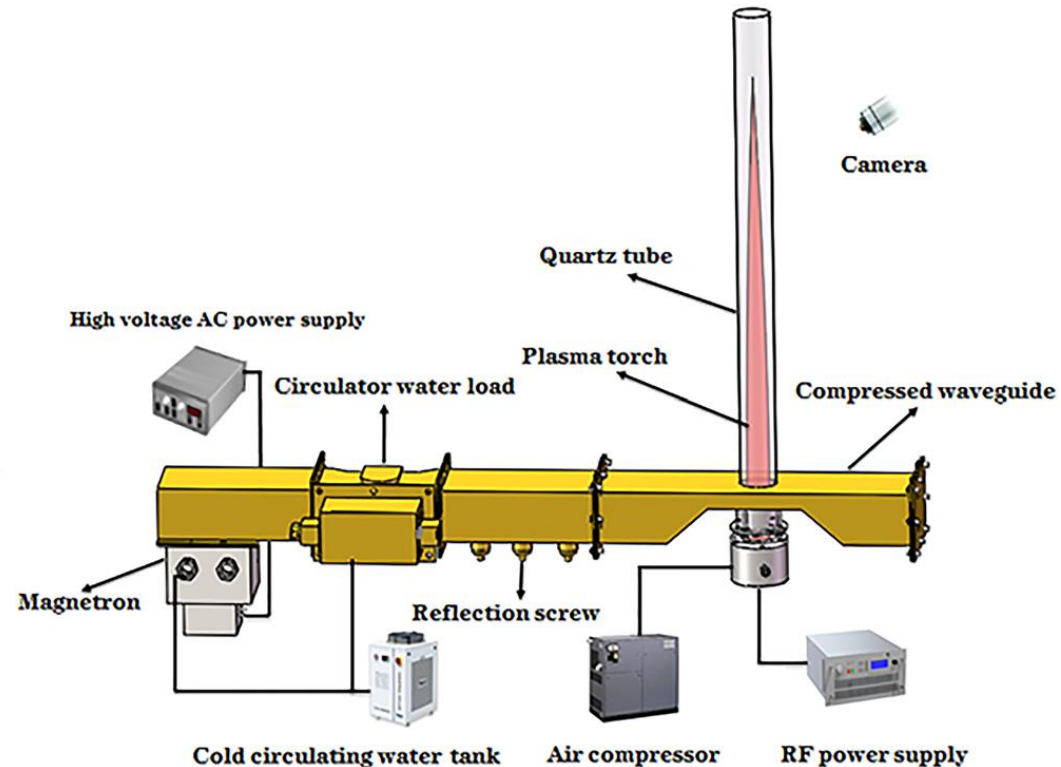
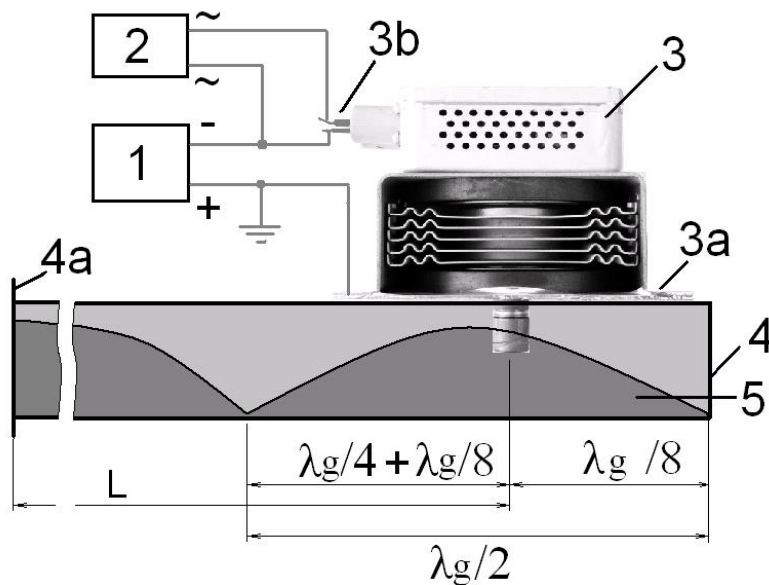


Magnetron schematic diagram

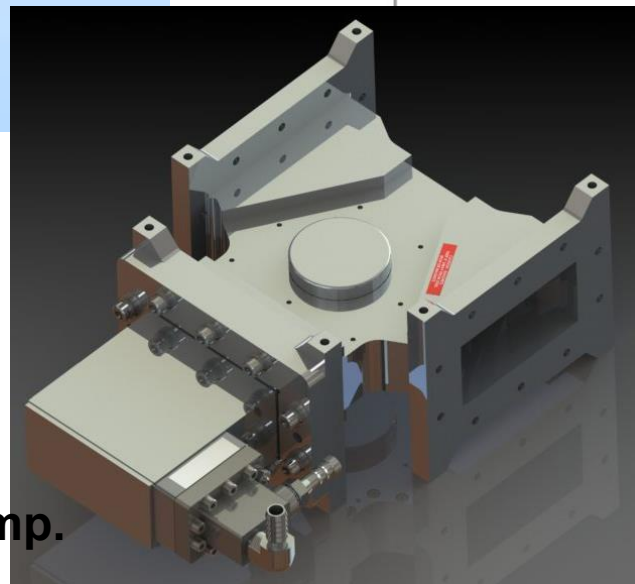
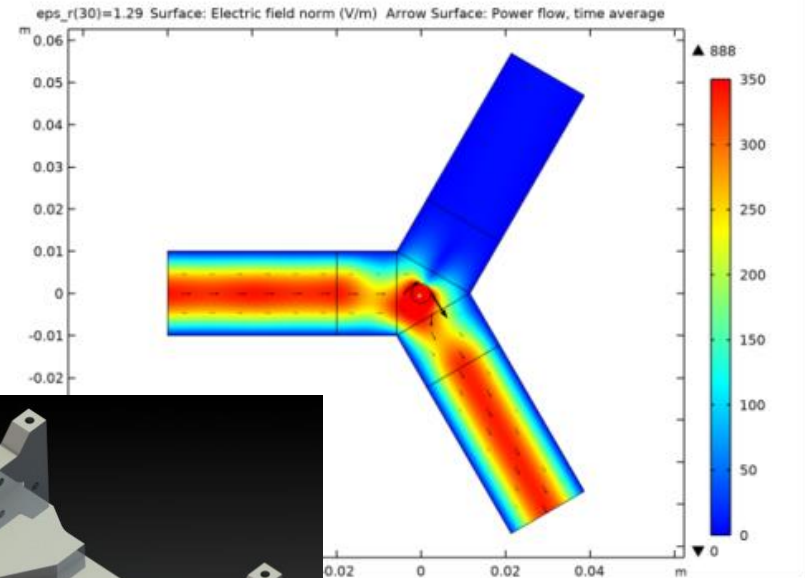
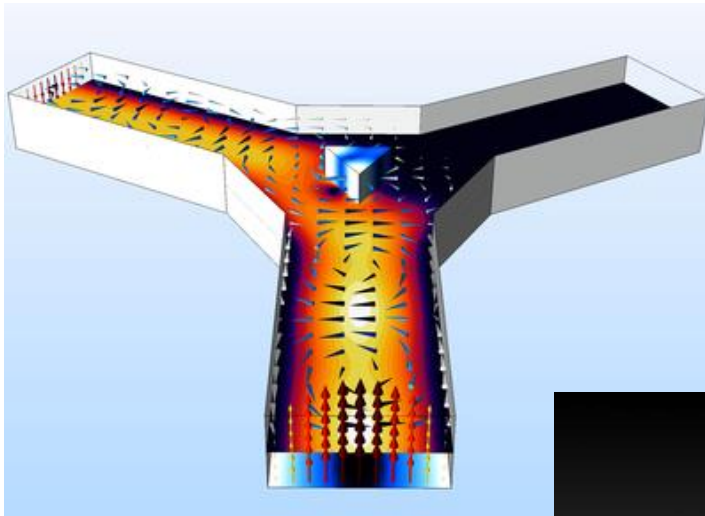


Microwave is generated.

The electrode of the microwave source is located at the location with the highest electric field



A 3-port circulator combining with a dump can be used as a isolator



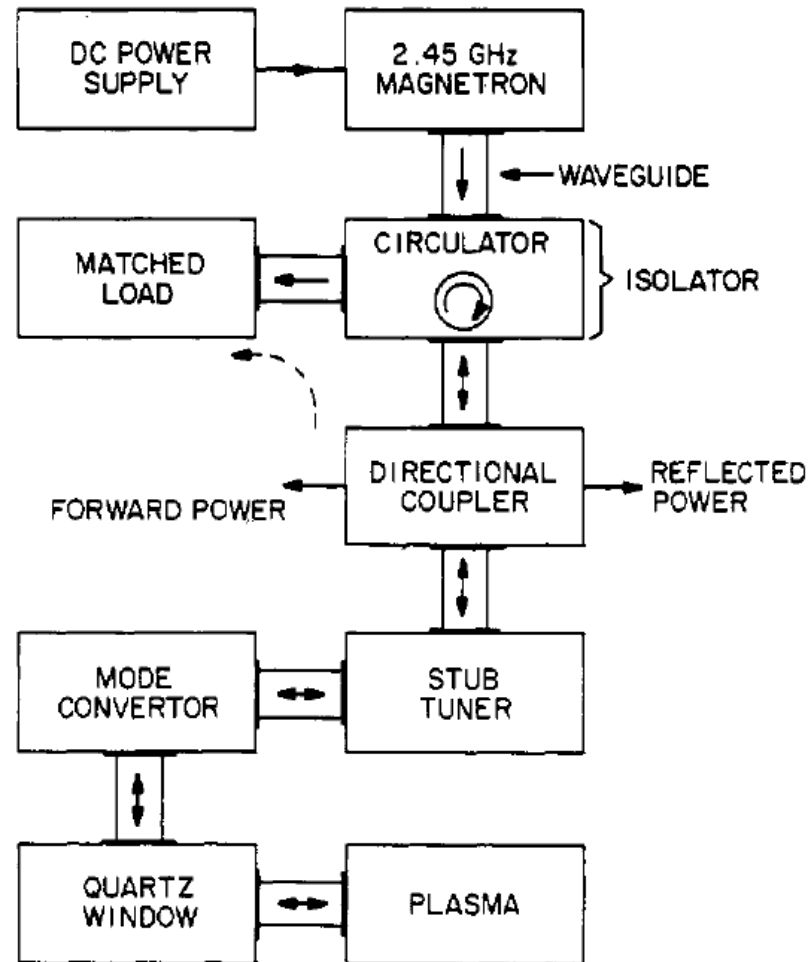
Water load as a dump.

<https://cn.comsol.com/model/impedance-matching-of-a-lossy-ferrite-3-port-circulator-10302>

<https://doc.comsol.com/6.0/doc/com.comsol.help.models.rf.circulator/circulator.html>

<https://ferriteinc.com/high-power-microwave-circulators-isolators/wr340-waveguide-s-band/>

Electron cyclotron resonance (ECR) microwave systems



microwave systems

Course Outline

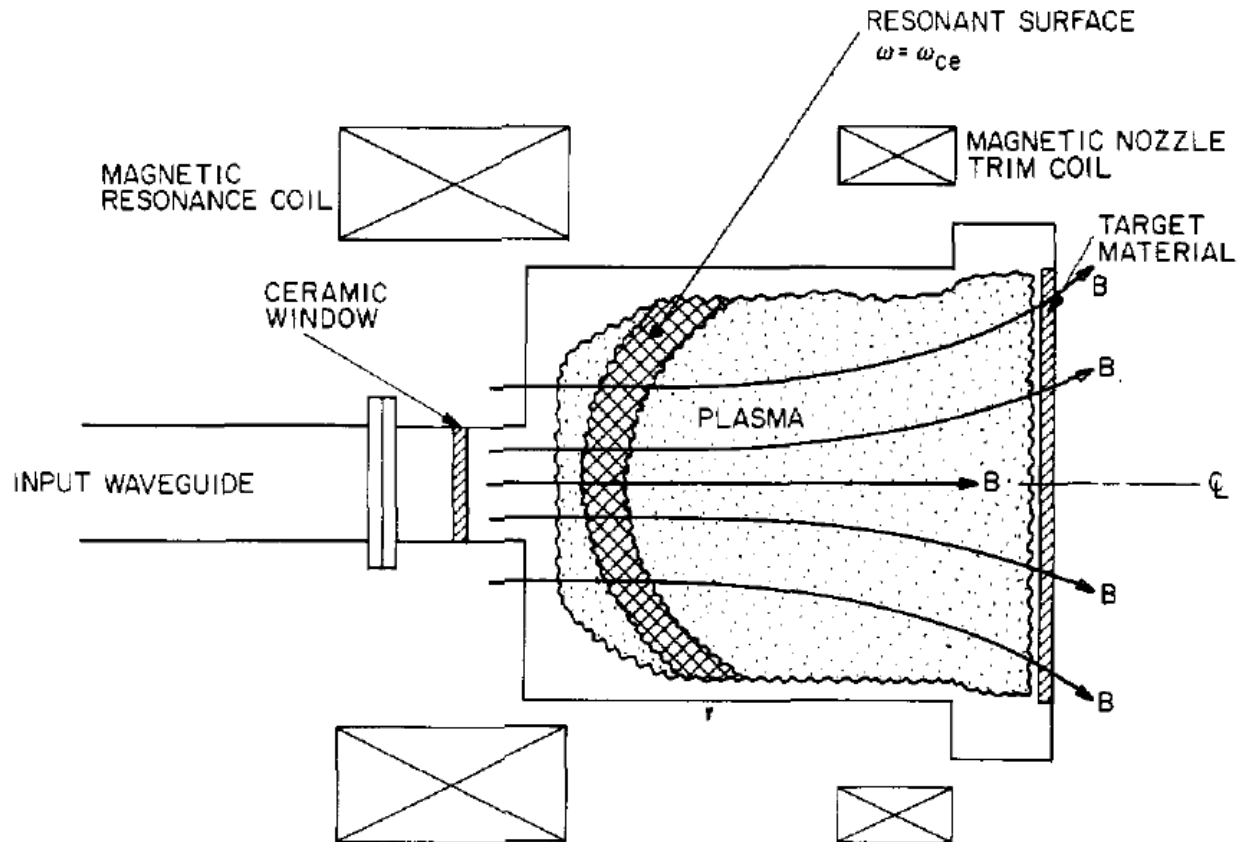


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Strong absorption occurs when the frequency matches the electron cyclotron frequency



- Electron cyclotron resonance (ECR) plasma reactor



Electron cyclotron frequency depends on magnetic field only



$$m_e \frac{d\vec{v}}{dt} = -\frac{e}{c} \vec{v} \times \vec{B}$$

- Assuming $\vec{B} = B\hat{z}$ and the electron oscillates in x-y plane

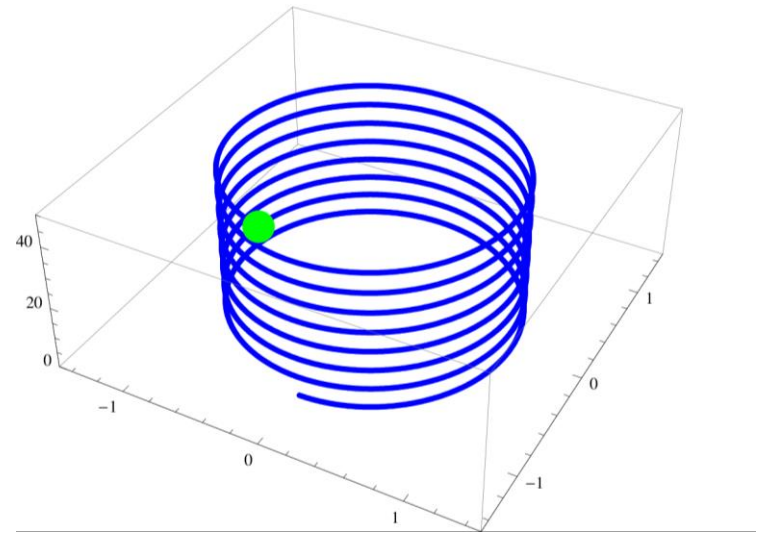
$$\begin{aligned} m_e \dot{v}_x &= -\frac{e}{c} B v_y & m_e \dot{v}_z &= 0 \\ m_e \dot{v}_y &= \frac{e}{c} B v_x \end{aligned}$$

$$\ddot{v}_x = -\frac{eB}{m_e c} \dot{v}_y = -\left(\frac{eB}{m_e c}\right)^2 v_x$$

$$\ddot{v}_y = -\frac{eB}{m_e c} \dot{v}_x = -\left(\frac{eB}{m_e c}\right)^2 v_y$$

- Therefore

$$\omega_{ce} = \frac{eB}{m_e c}$$



Electrons keep getting accelerated when a electric field rotates in electron's gyrofrequency



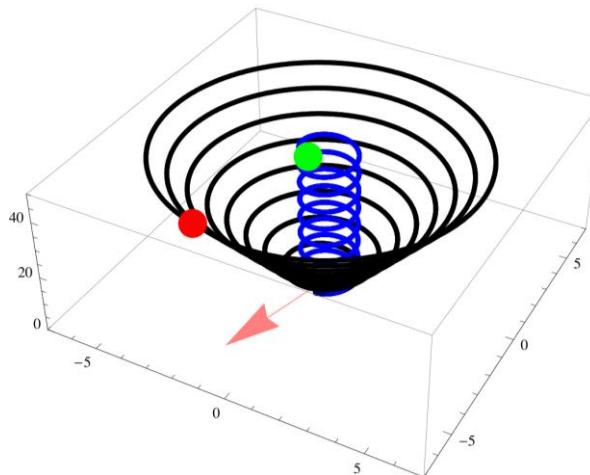
$$m_e \frac{d\vec{v}}{dt} = -\frac{e}{c} \vec{v} \times \vec{B} - e \vec{E} \quad \vec{B} = B_0 \hat{z} \quad \vec{E} = E_0 [\hat{x} \cos(\omega t) + \hat{y} \sin(\omega t)]$$

$$m_e \dot{v}_x = -\frac{e}{c} B v_y + E_0 \cos(\omega t) \quad m_e \dot{v}_y = \frac{e}{c} B v_x + E_0 \sin(\omega t) \quad m_e \dot{v}_z = 0$$

$$\ddot{v}_x = -\frac{eB}{m_e c} \dot{v}_y - \frac{E_0}{m_e} \omega \sin(\omega t) = -\omega_{ce}^2 v_x - \frac{E_0}{m_e} (\omega_{ce} + \omega) \sin(\omega t)$$

$$\ddot{v}_y = -\frac{eB}{m_e c} \dot{v}_x + \frac{E_0}{m_e} \omega \cos(\omega t) = -\omega_{ce}^2 v_y + \frac{E_0}{m_e} (\omega_{ce} + \omega) \cos(\omega t)$$

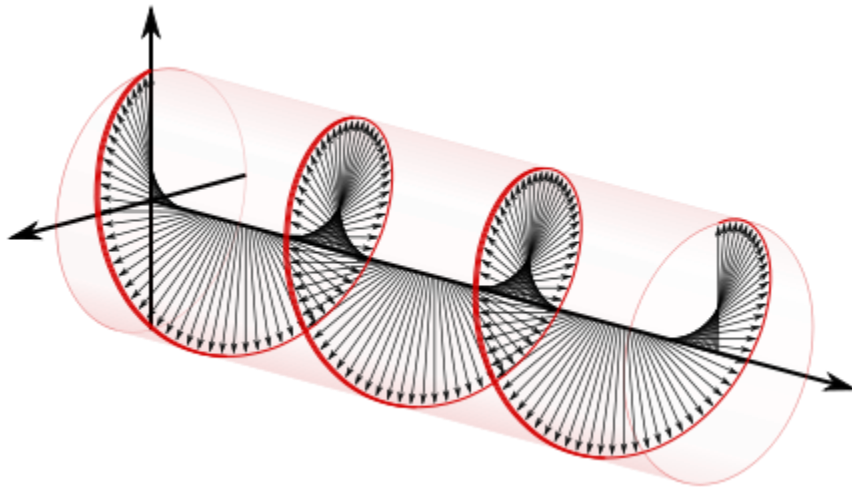
$$\omega_{ce} = \frac{eB}{m_e c}$$



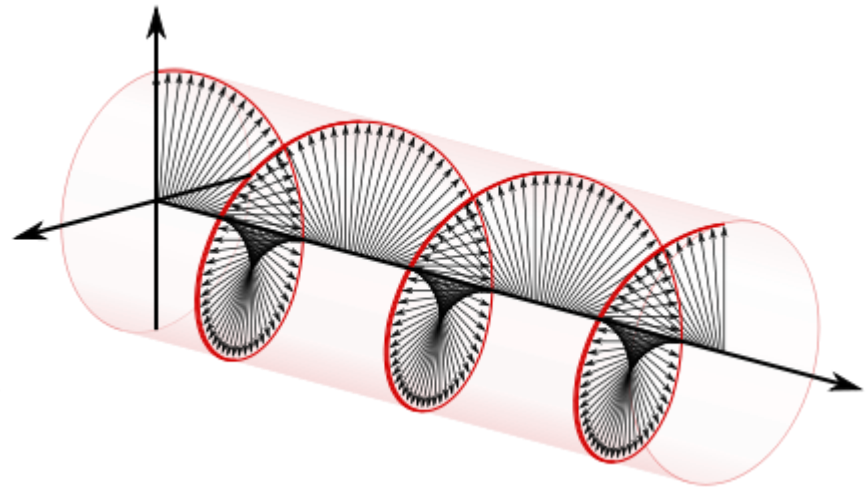
Electric field in a circular polarized electromagnetic wave keeps rotating as the wave propagates



- Right-handed polarization



- Left-handed polarization

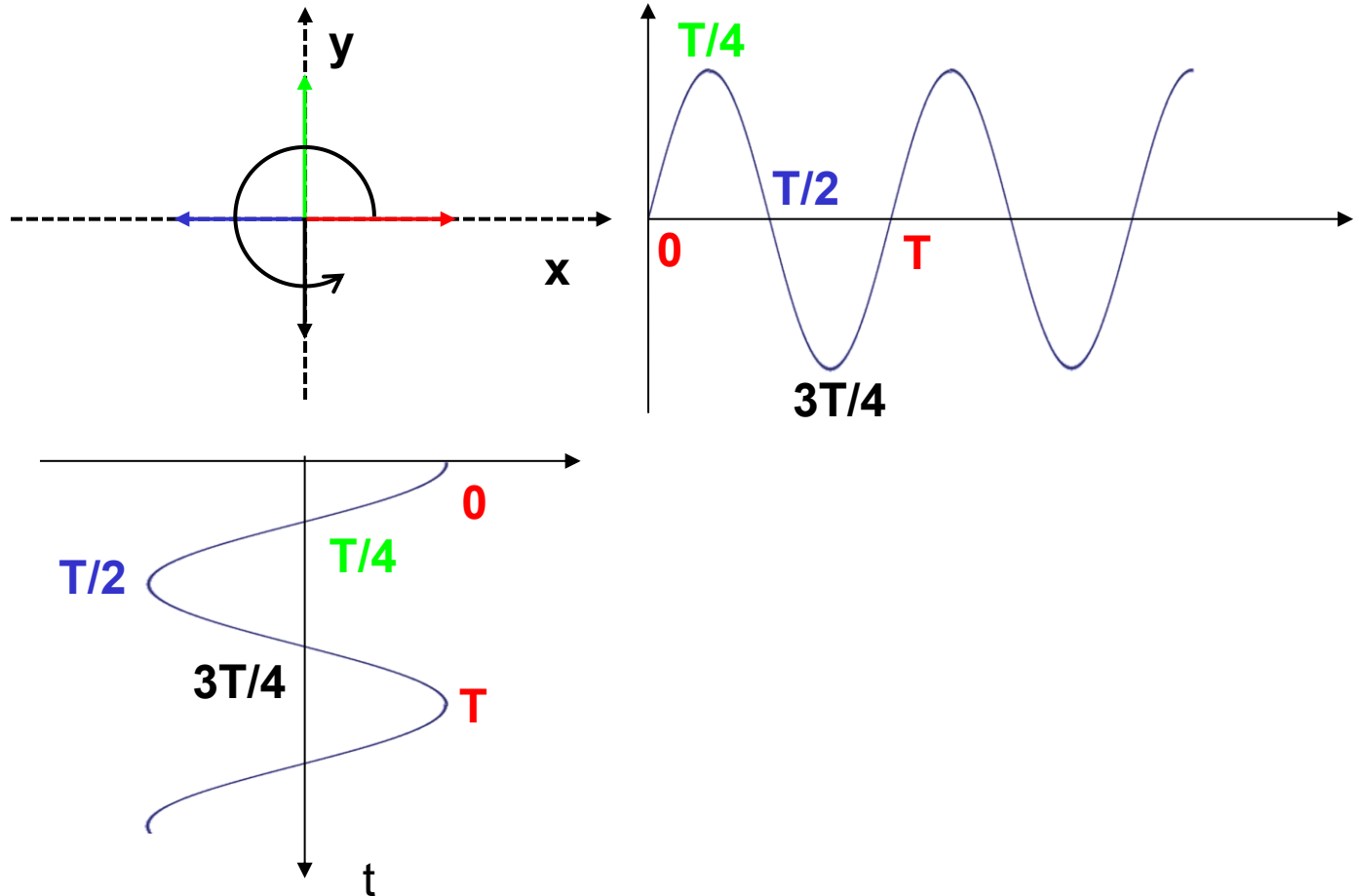


Electric field rotates in a circular polarization



$$E_x = E_0 \exp(-i\omega t)$$

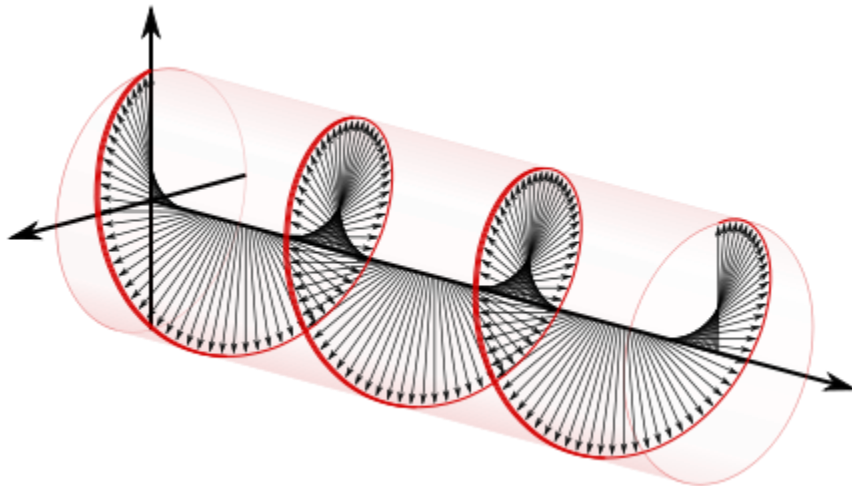
$$E_y = \pm i E_x = i E_0 \exp(-i\omega t) = E_0 \exp\left(\pm i \frac{\pi}{2}\right) \exp(-i\omega t) = E_0 \exp\left[-i\left(\omega t \pm \frac{\pi}{2}\right)\right]$$



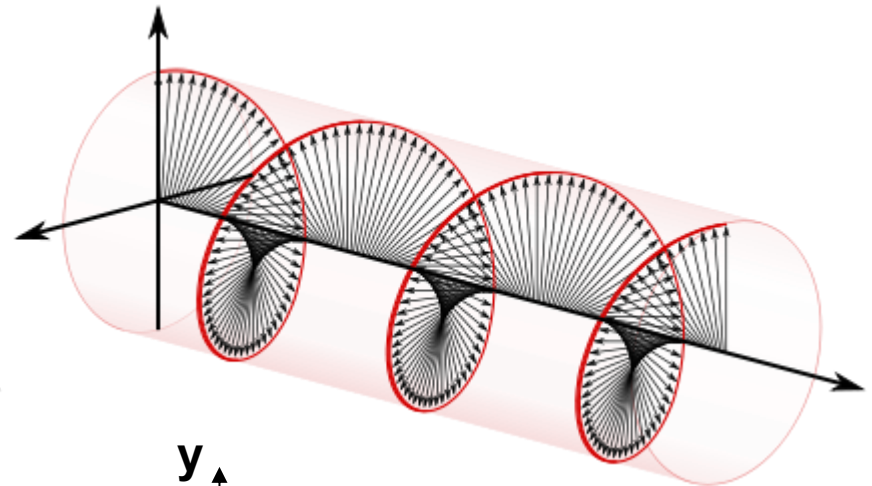
A linear polarized wave can be decomposed by a left-handed and a right-handed polarized wave



- Right-handed polarization

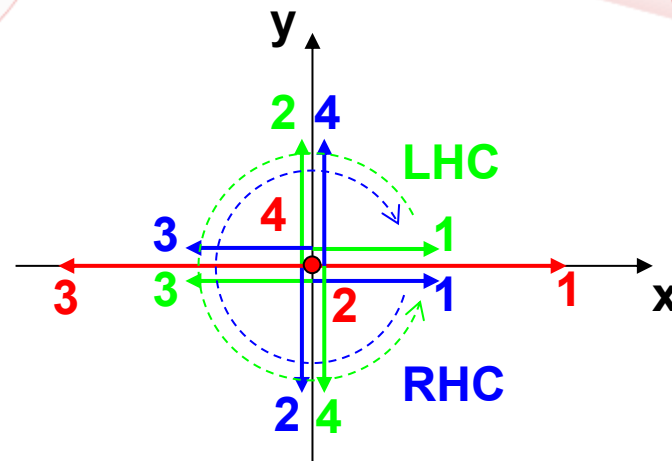


- Left-handed polarization



$$\vec{E} = E_0 \hat{x} = \frac{E_0}{2} [(\hat{x} + i\hat{y}) + (\hat{x} - i\hat{y})]$$

RHC LHC



Only right-handed polarization can resonance with electron's gyromotion

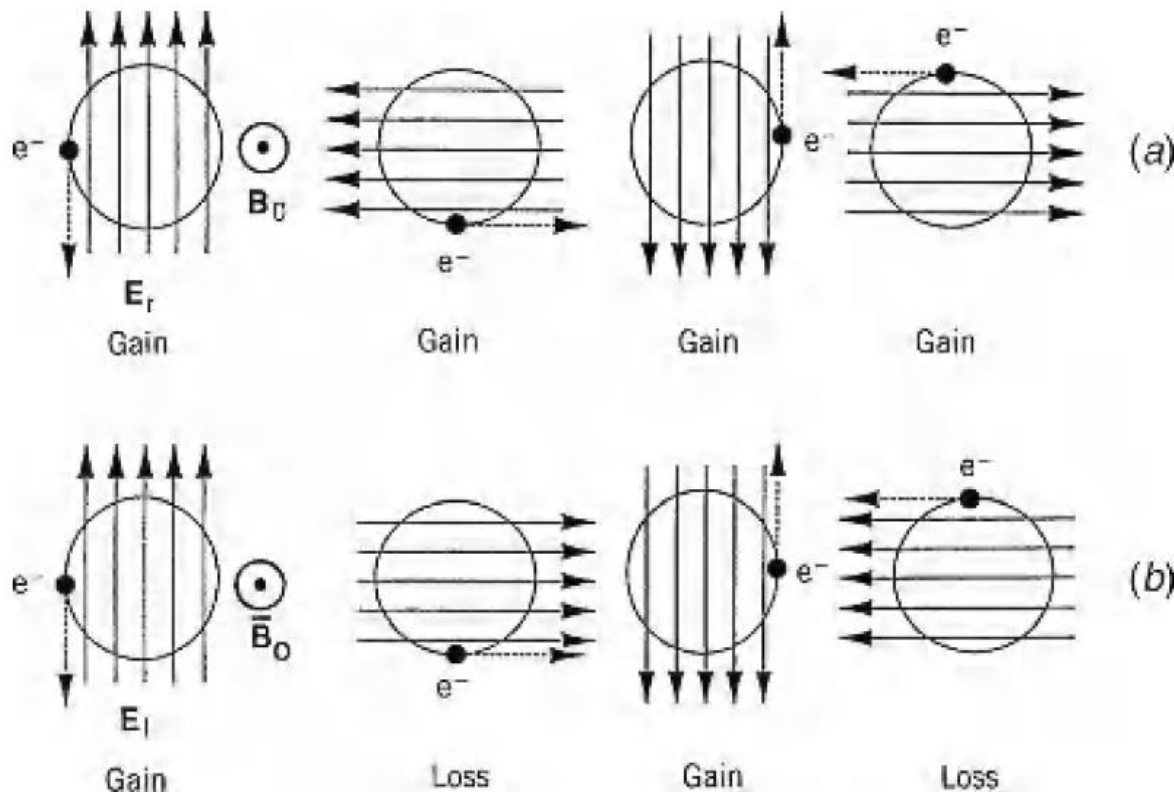
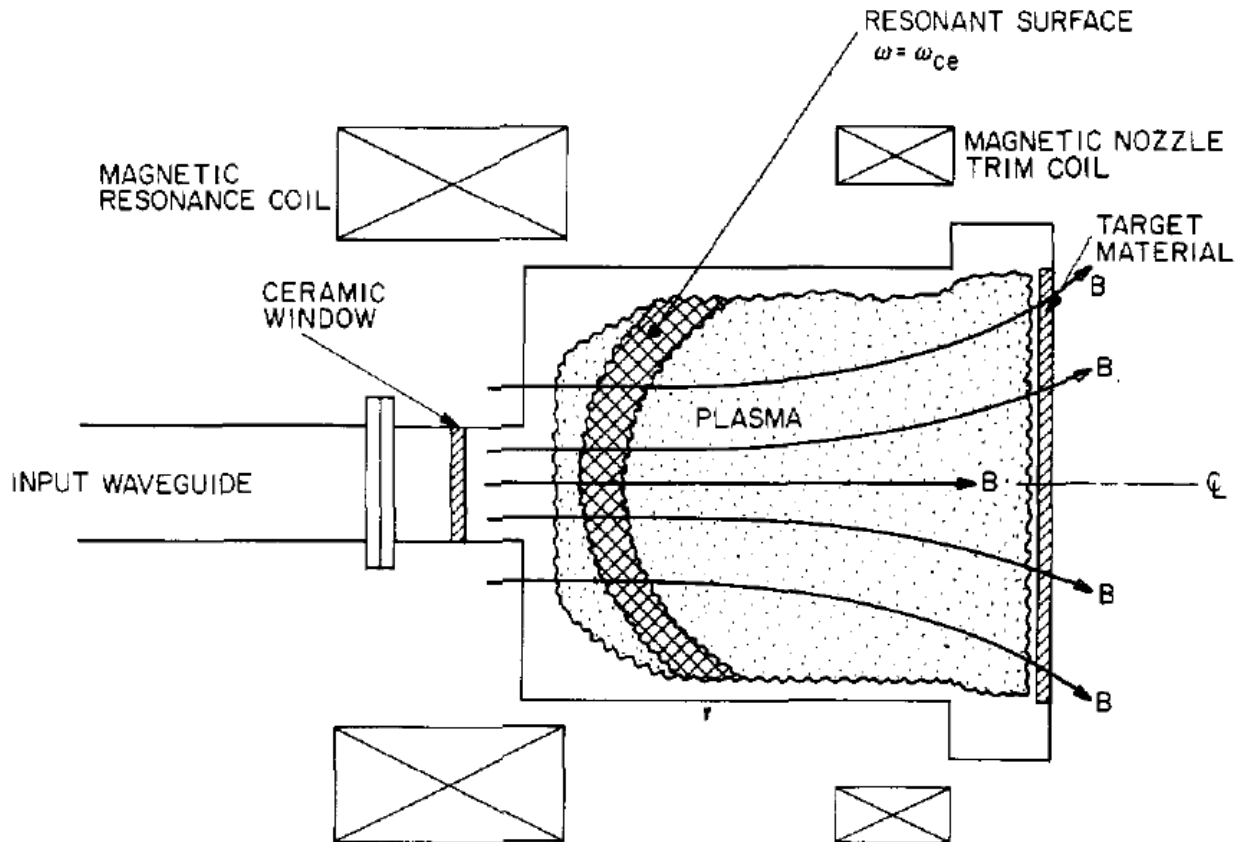


FIGURE 13.5. Basic principle of ECR heating: (a) continuous energy gain for right-hand polarization; (b) oscillating energy for left-hand polarization (after Lieberman and Gottscho, 1994).

Strong absorption occurs when the frequency matches the electron cyclotron frequency



- Electron cyclotron resonance (ECR) plasma reactor

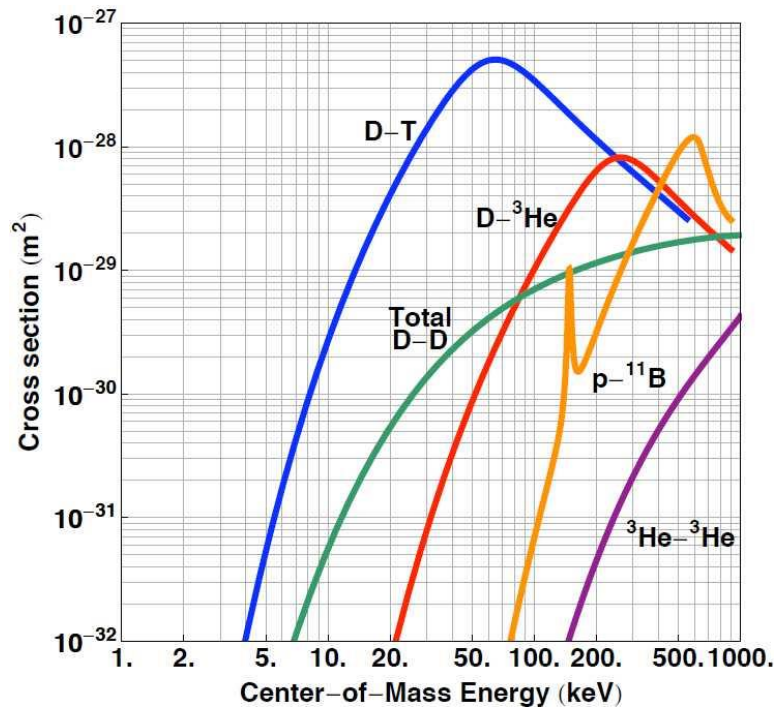


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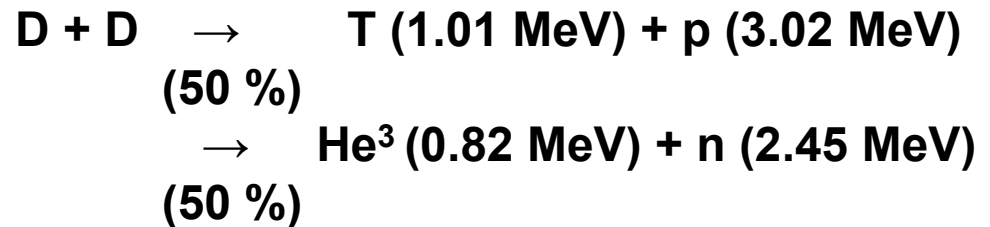
Two nuclear fusion reactions will be observed in experiments



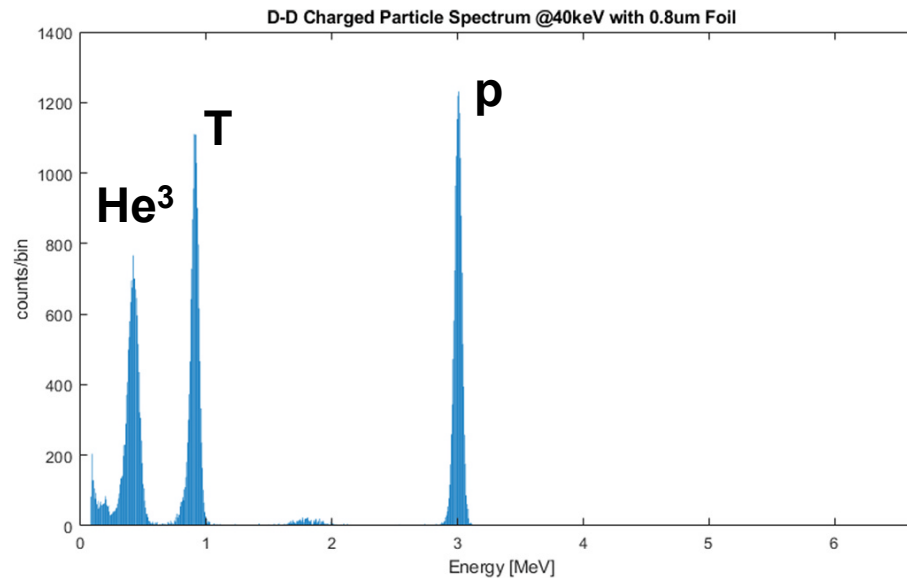
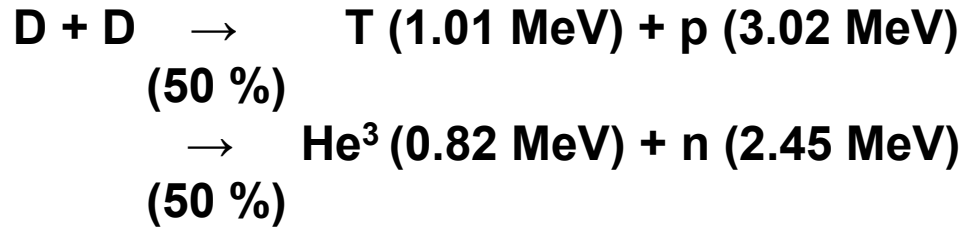
- Local experiment:



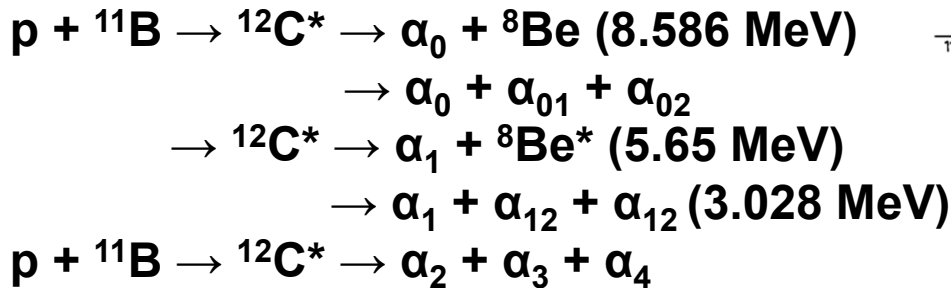
- Remote experiment:



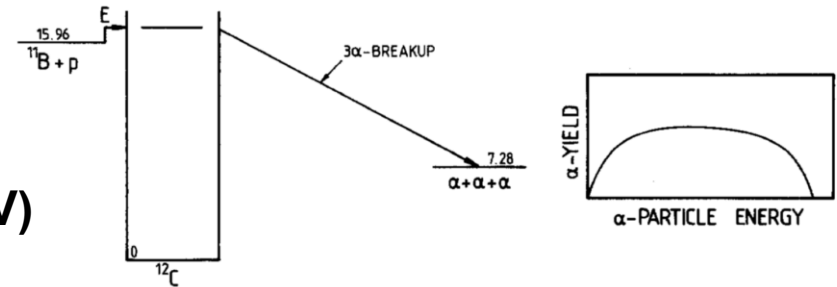
Monoenergetic particles will be generated in D-D fusion



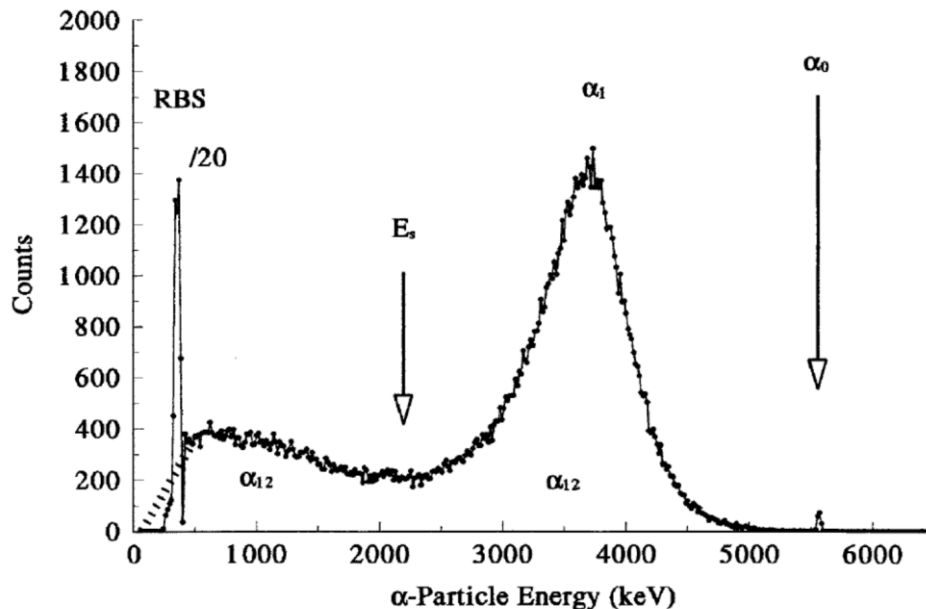
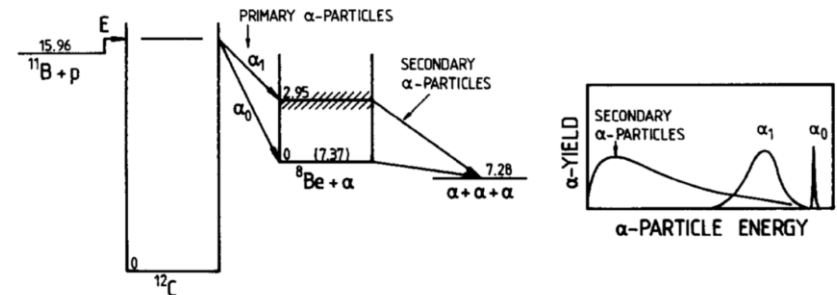
Alphas with wide spectrum will be generated



(a) DIRECT DECAY



(b) SEQUENTIAL DECAY



**RBS: Rutherford
backscattering
spectrometry**

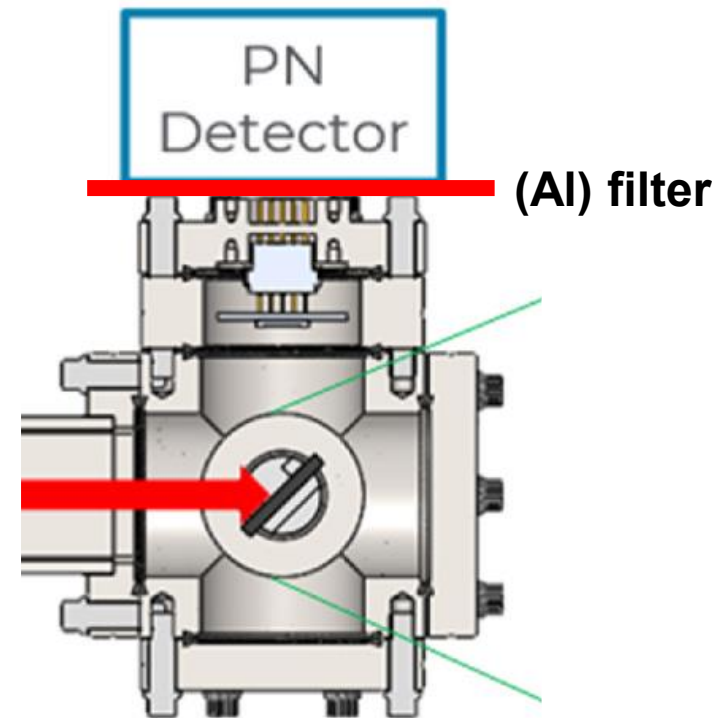
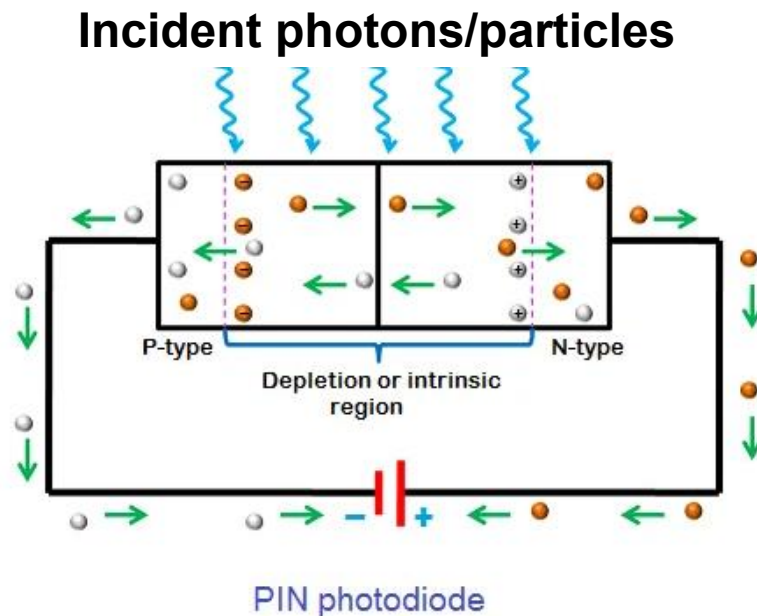
H. W. Becker, etc, Z. Phys. A – Atomic Nuclei 327, 341 (1987)
 Jiarui Liu, etc., Nucl. Instr. And Meth. B 190, 107 (2002)
 F. A. Geser and M. Valente, Radiat. Phys. Chem 167, 108224 (2020)

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A filter is used to block the photon so that the detector only sees alphas



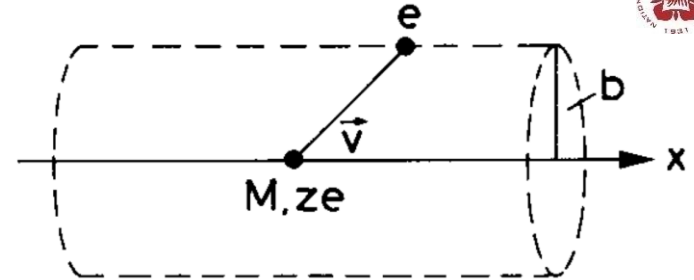
Energetic charged particles losses most of its energy right before it stops



- Momentum transfer:

$$\Delta p_{\perp} = \int F_{\perp} dt = \int F_{\perp} \frac{dt}{dx} dx = \int F_{\perp} \frac{dx}{v}$$

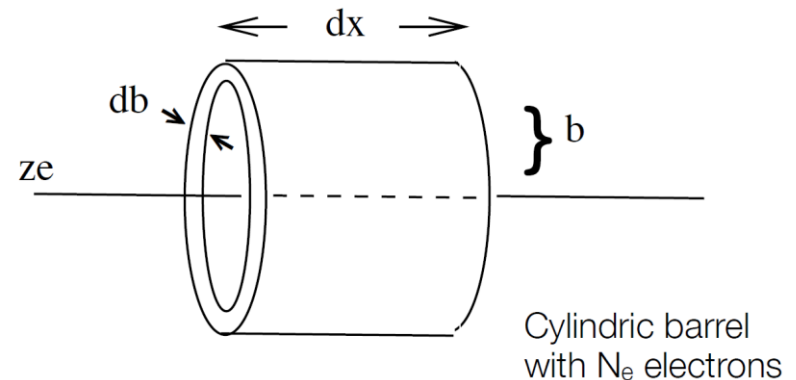
$$= \int_{-\infty}^{\infty} \frac{ze^2}{(x^2 + b^2)} \frac{b}{\sqrt{x^2 + b^2}} \frac{1}{v} dx = \frac{ze^2 b}{v} \left[\frac{x}{b^2 \sqrt{x^2 + b^2}} \right]_{-\infty}^{\infty} = \frac{2ze^2}{bv}$$



$\Delta p_{\parallel}$: averages to zero due to symmetry

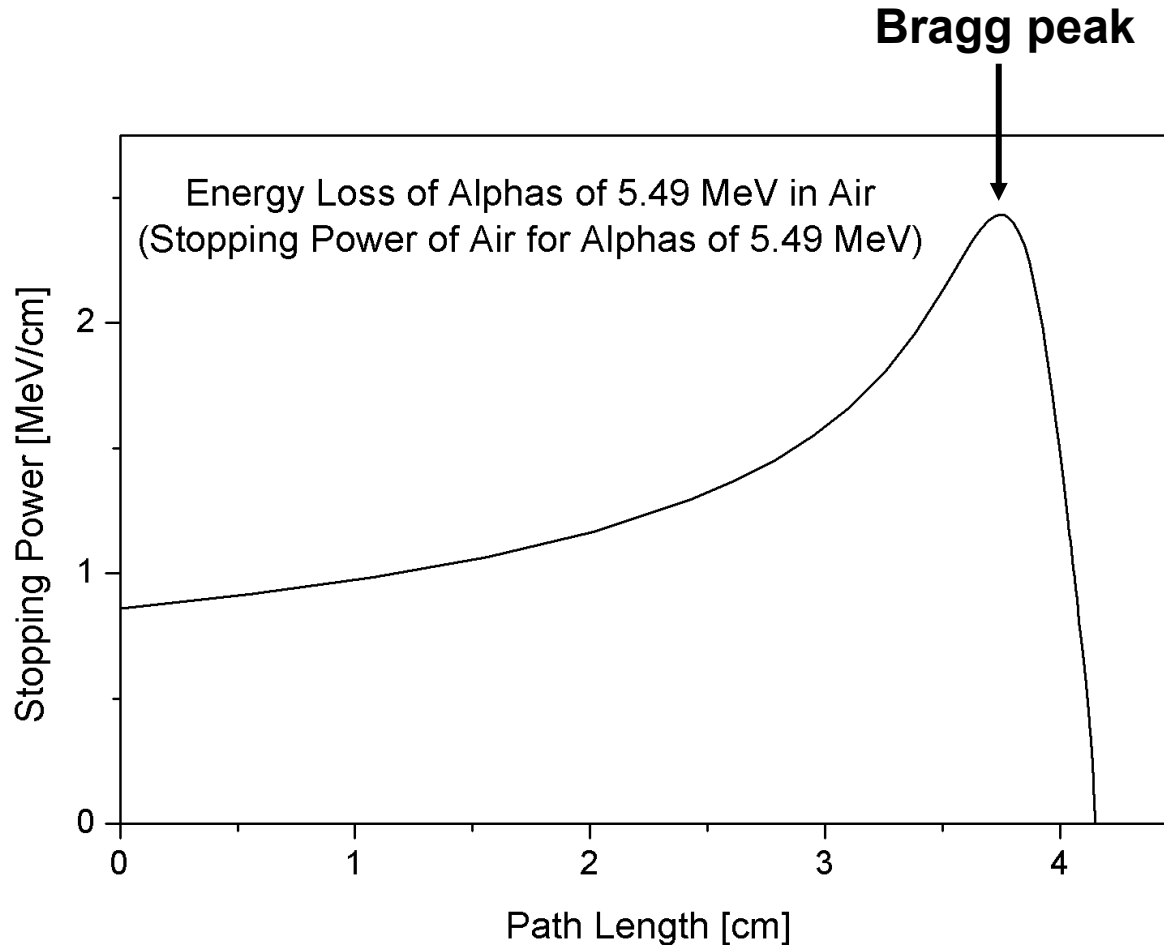
$$\Delta E(b) = \frac{\Delta p^2}{2m_e} \quad N_e = n(2\pi b)dbdx$$

$$-dE(b) = \frac{\Delta p^2}{2m_e} 2\pi n b db dx$$



$$-\frac{dE}{dx} = \frac{4\pi n z^2 e^4}{m_e v^2} \int_{b_{\min}}^{b_{\max}} \frac{db}{b} = \frac{4\pi n z^2 e^4}{m_e v^2} \ln \left(\frac{b_{\max}}{b_{\min}} \right) \propto \frac{1}{E}$$

A particle loses most of its energy right before it stops



There are two suggested website for getting the information of proton stopping power in different materials



<http://www.nist.gov/pml/data/star/>

<http://www.srim.org/>

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NIST Home > PML > Physical Reference Data > Stopping-Power & Range Tables: e-, p+, Helium Ions

NISTIR 4999 | [Version History](#) | [Disclaimer](#)

Stopping-Power and Range Tables for Electrons, Protons, and Helium Ions

M.J. Berger, J.S. Coursey, M.A. Zucker and J. Chang
(NIST, Physical Measurement Laboratory)



Abstract:
The databases ESTAR, PSTAR, and ASTAR calculate stopping-power and range tables for electrons, protons, or helium ions, according to methods described in ICRU Reports 37 and 49. Stopping-power and range tables can be calculated for electrons in any user-specified material and for protons and helium ions in 74 materials.

Contents:

1. Introduction
2. ESTAR: Stopping Powers and Ranges for Electrons
3. PSTAR and ASTAR: for Protons and Helium Ions (alpha particles)

[References](#)

[Appendix: Significance of Calculated Quantities](#)

Access the Data:

1. Electrons
2. Protons

請選擇語言

由「Google 翻譯」技術提供



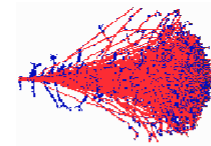
© Creations/2010 Shutterstock.com

Access the Data

[Electrons](#) | [Protons](#) | [Helium Ions](#)

NIST Standard Reference Database 124
Rate our products and services.
Online: October 1998 - **Last update:** August 2005

Contact
Stanley Seltzer



SRIM Textbook

Software	Science
SRIM / TRIM Introduction	Historical Review
Download SRIM-2013	Details of SRIM-2013
SRIM Install Problems	Experimental Data Plots
SRIM Tutorials	Stopping of Ions in Matter
Download TRIM Manual Part-1, Part-2	Stopping in Compounds
Scientific Citations of Experimental Data	High Energy Stopping

The thickness of a filter can be decided from the range data from NIST website



COPPER

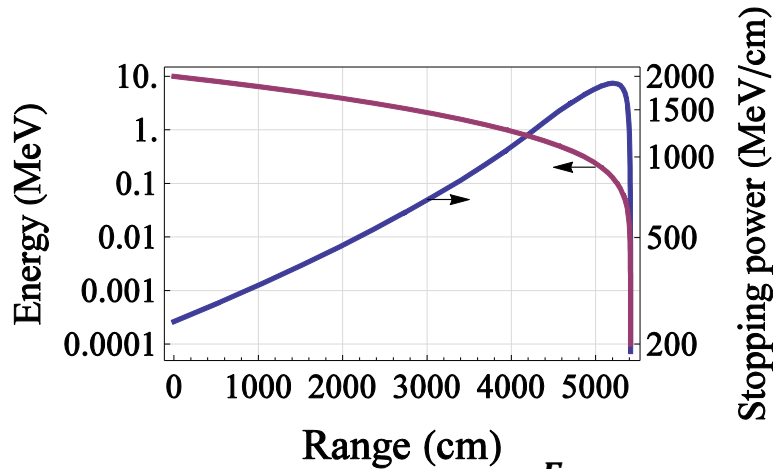
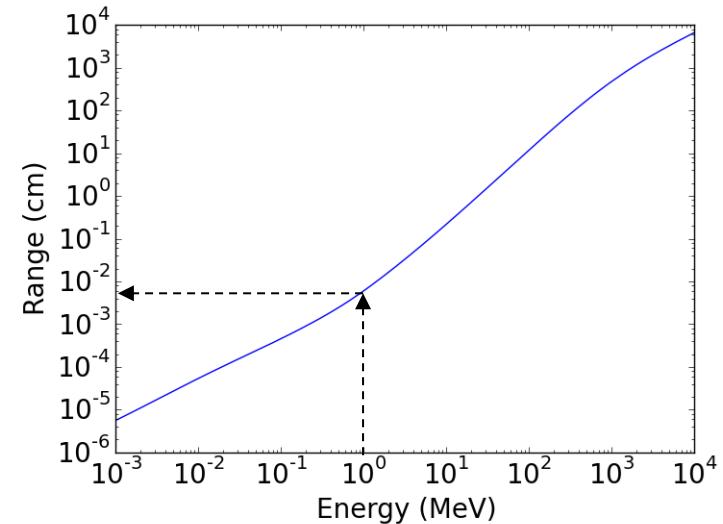
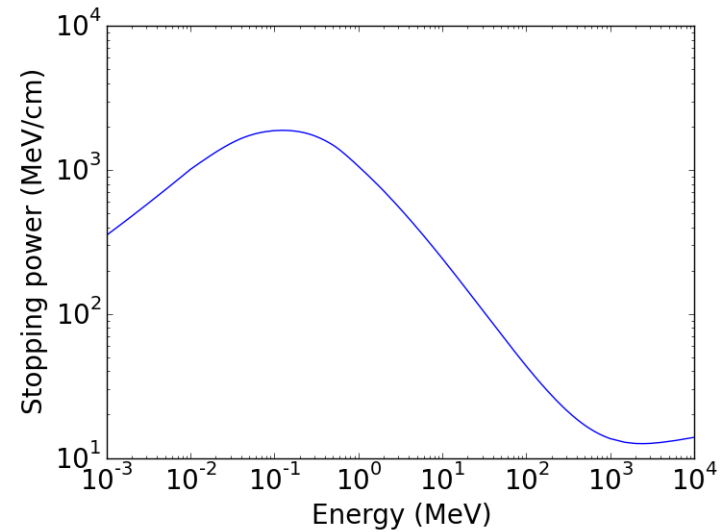
To download data in spreadsheet (array) form, choose a delimiter and use the checkboxes in the table heading. After downloading, save the output by using your browser's Save As feature.

Delimiter:

- ☒ space
- ☐ (vertical bar)
- ☐ tab (some browsers may use spaces instead)
- ☐ newline

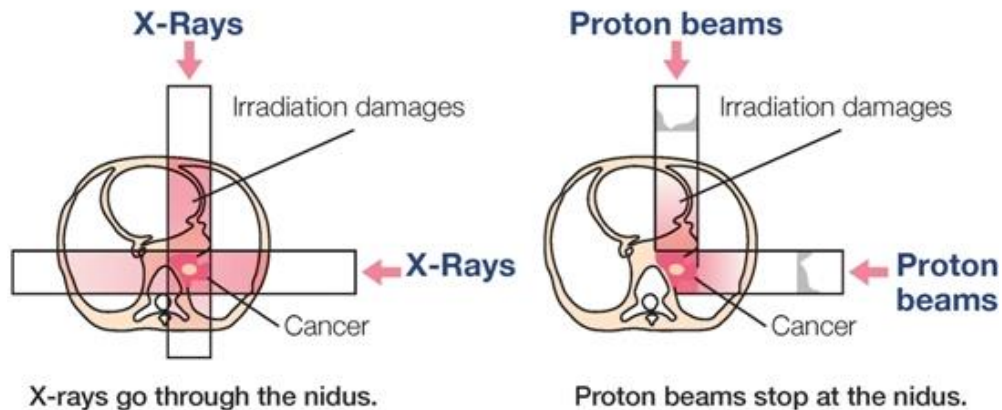
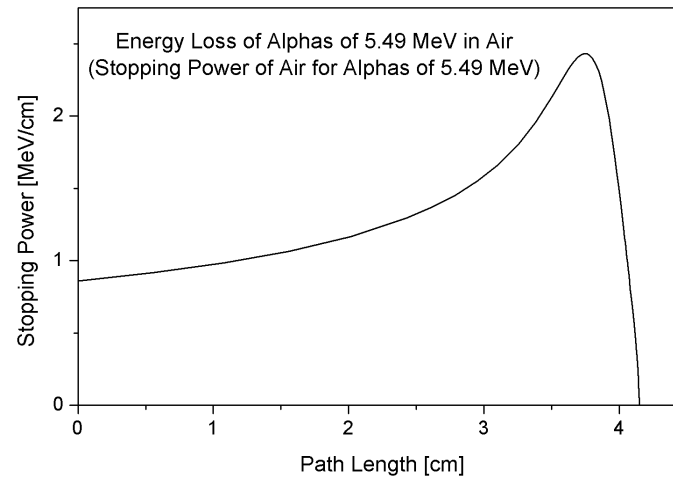
Download data [Reset]

(required) Kinetic Energy (MeV)	Stopping Power (MeV cm ² /g)			Range		
	Electronic	Nuclear	Total	CSDA (g/cm ²)	Projected (g/cm ²)	Detour Factor Projected / CSDA
1.000E-03	3.490E+01	4.408E+00	3.931E+01	4.116E-05	5.620E-06	0.1365
1.500E-03	4.274E+01	4.231E+00	4.697E+01	5.267E-05	8.301E-06	0.1576
2.000E-03	4.935E+01	4.049E+00	5.340E+01	6.263E-05	1.101E-05	0.1759
2.500E-03	5.518E+01	3.876E+00	5.906E+01	7.152E-05	1.374E-05	0.1921
3.000E-03	6.045E+01	3.718E+00	6.416E+01	7.964E-05	1.647E-05	0.2068
4.000E-03	6.980E+01	3.440E+00	7.324E+01	9.419E-05	2.194E-05	0.2329
5.000E-03	7.804E+01	3.207E+00	8.124E+01	1.071E-04	2.739E-05	0.2556
6.000E-03	8.548E+01	3.010E+00	8.849E+01	1.189E-04	3.280E-05	0.2758
7.000E-03	9.233E+01	2.840E+00	9.517E+01	1.298E-04	3.817E-05	0.2940
8.000E-03	9.871E+01	2.692E+00	1.014E+02	1.400E-04	4.347E-05	0.3106
9.000E-03	1.047E+02	2.561E+00	1.073E+02	1.496E-04	4.872E-05	0.3258
1.000E-02	1.104E+02	2.445E+00	1.128E+02	1.587E-04	5.391E-05	0.3398



$$\frac{dE}{dx} = f(E) \Rightarrow x = \int_{E_i}^{E_f} \frac{dE}{f(E)}$$

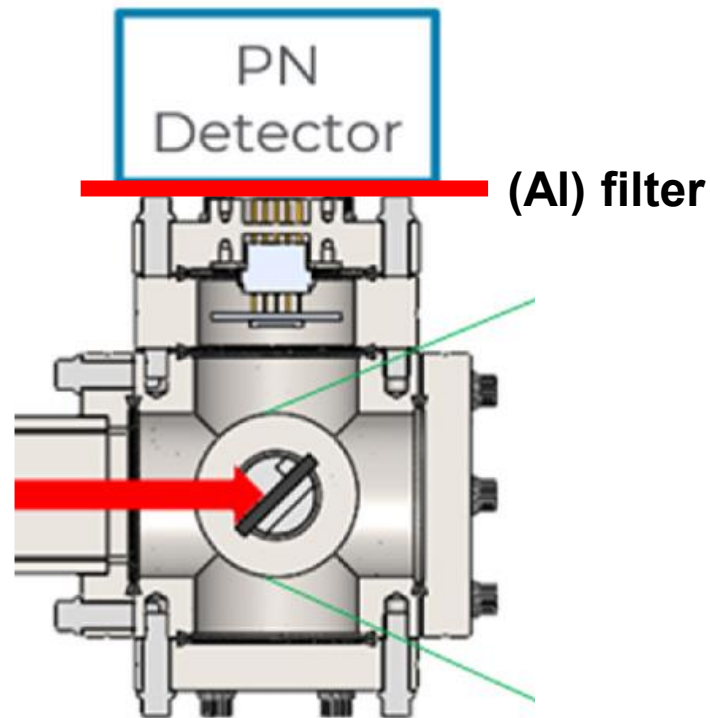
Proton therapy takes the advantage of using Bragg peak



Alpha particles loss energies when passing through the (Al) filter



- A filter is used to block the photon so that the detector only sees alphas.
- Alphas in $p\text{-}^{11}\text{B}$, which are mostly in the range of 3–5 MeV, would be attenuated by ~130–180 keV in 0.8 μm of Al.



Course Outline

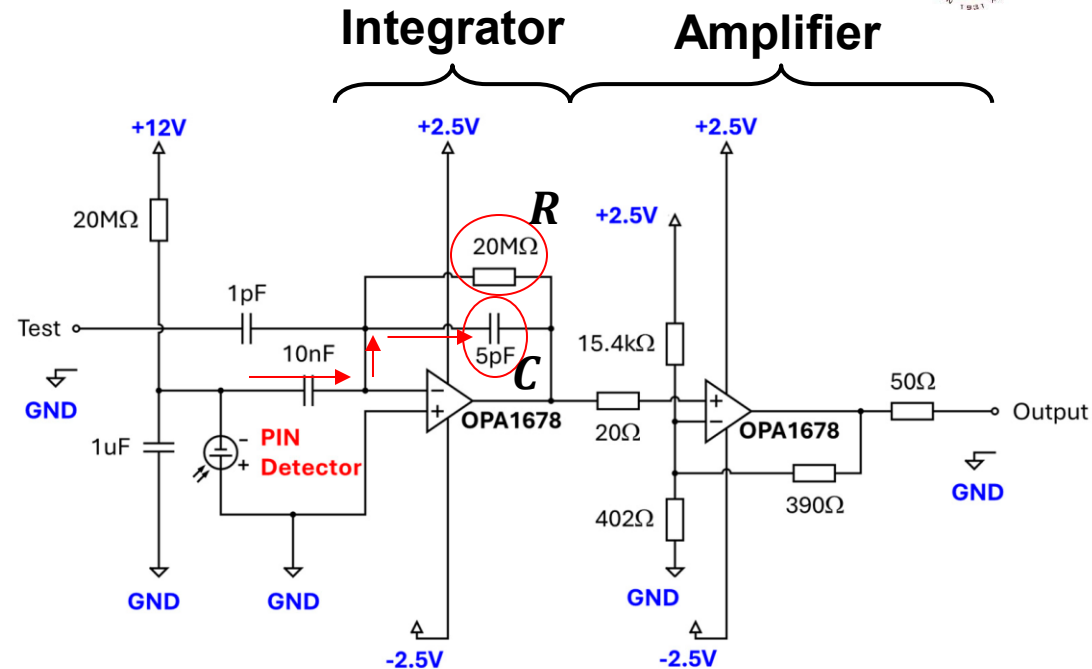
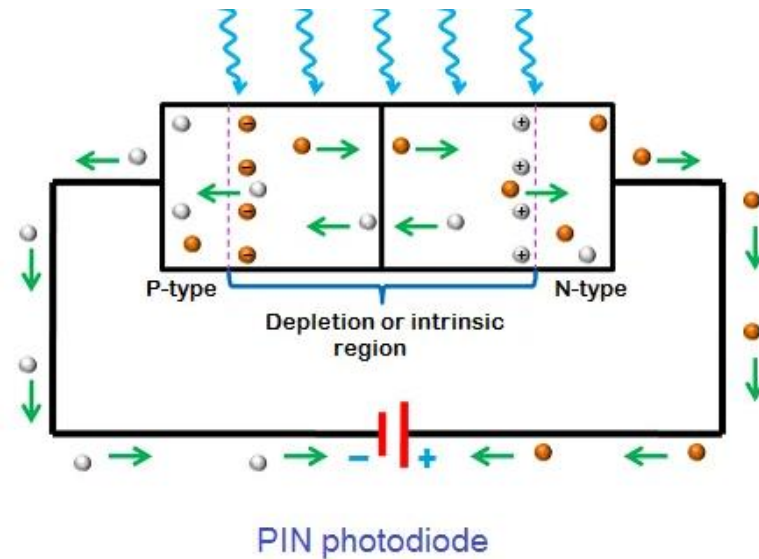


- Formosa Integrated Research Spherical Tokamak (FIRST), the first Tokamak in Taiwan
- **Physics used in Alpha-E**
 - Microwave
 - ECRH
 - Spectrum of fusion products
 - Stopping power of the filter
 - **Data analysis using Pulse Shaping**
- Final report

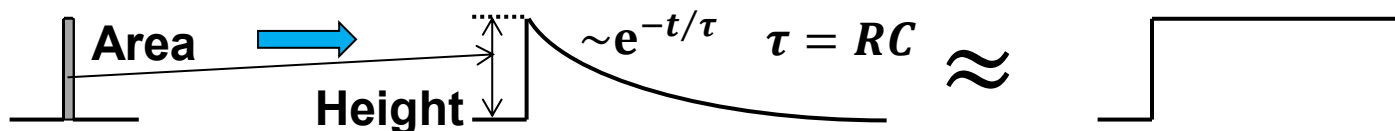
A pulse signal is converted to a voltage signal using an integrator



Incident photons/particles



- The 5-pF capacitor is charged by the current from the PIN detector.
- The 5-pF capacitor is discharged by the 20-MΩ resistor.



Pulse Shaping is commonly used in nuclear and particle physics electronics



- Improve Signal-to-noise ratio S/N
 - Restrict bandwidth to match measurement time.

=> Increase pulse width.

- Improve pulse-pair resolution

=> Decrease pulse width.

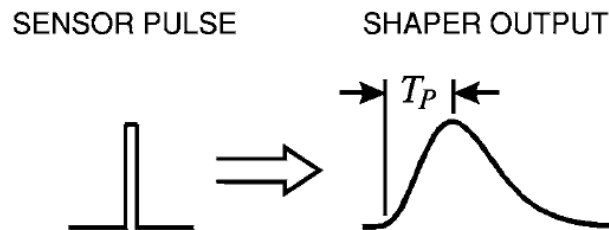


FIGURE 13. A pulse shaper transforms a short sensor pulse into a longer pulse with a rounded cusp and peaking time T_P .

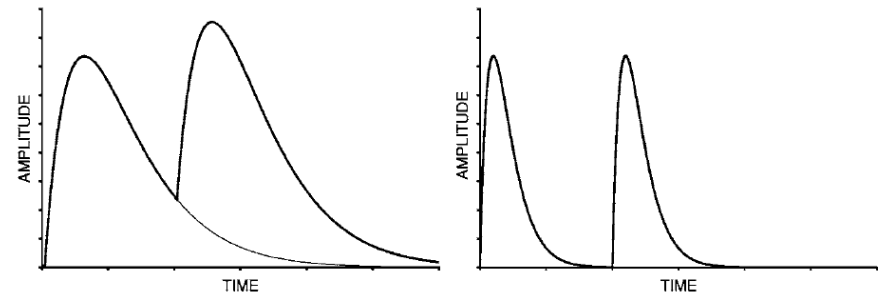
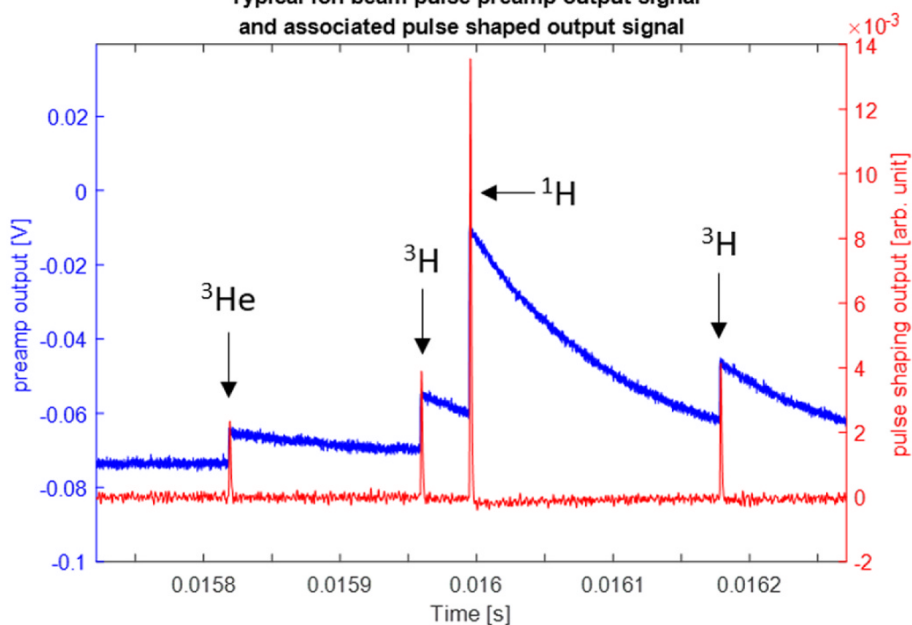


FIGURE 14. Amplitude pileup when two successive pulses overlap (left). Reducing the shaping time allows the first pulse to return to the baseline before the second arrives.

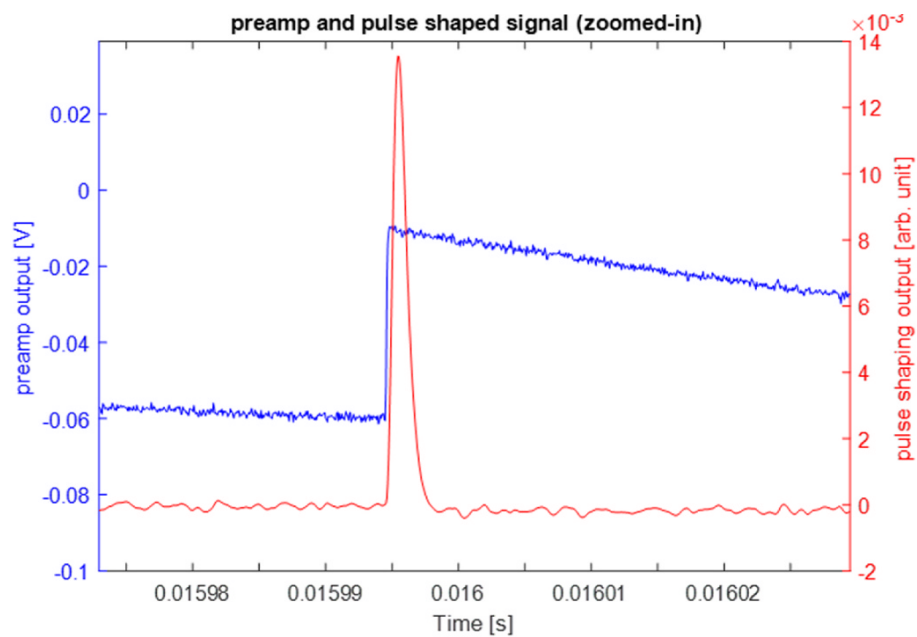
Expected data profile



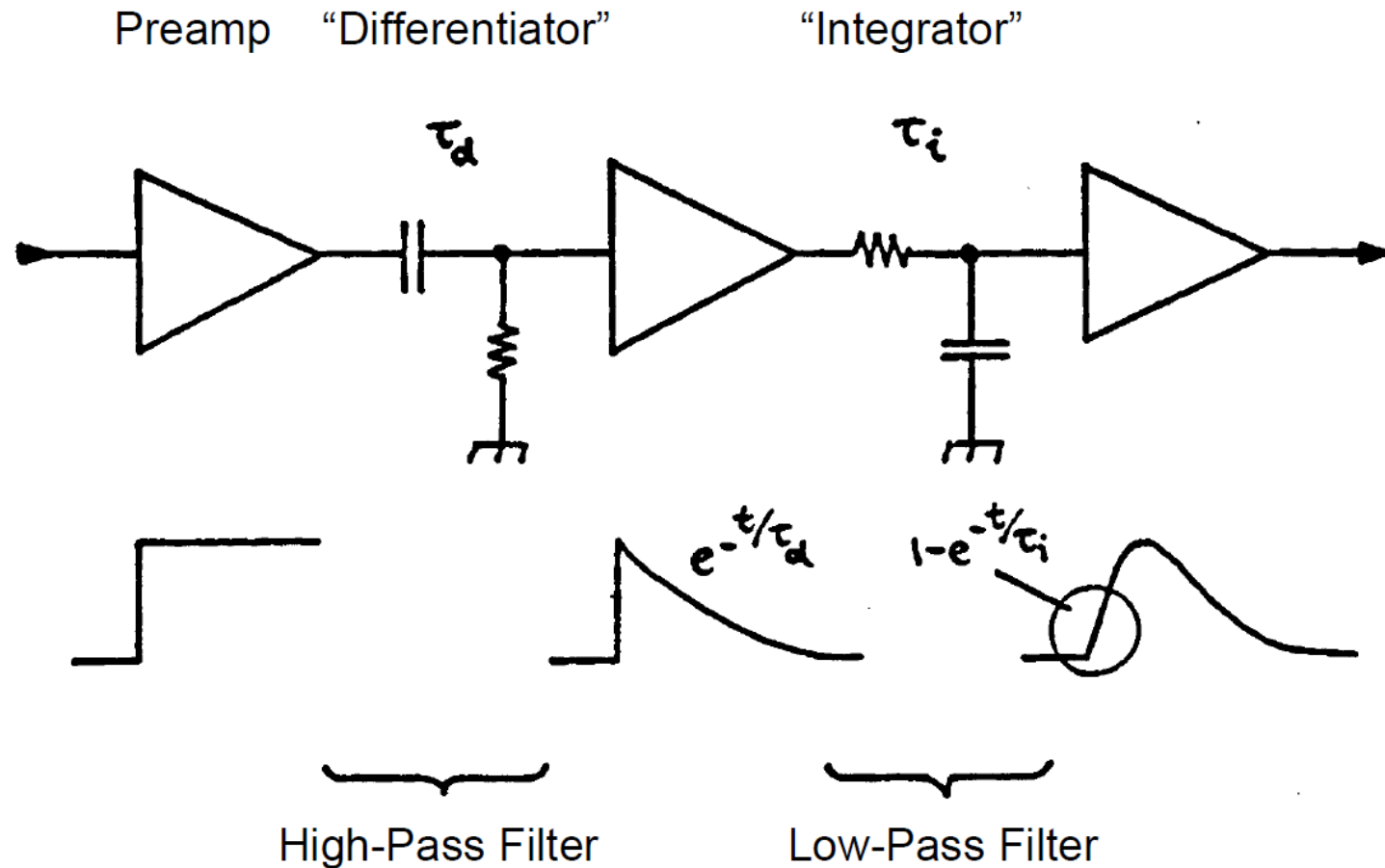
Typical ion beam pulse preamp output signal
and associated pulse shaped output signal



preamp and pulse shaped signal (zoomed-in)



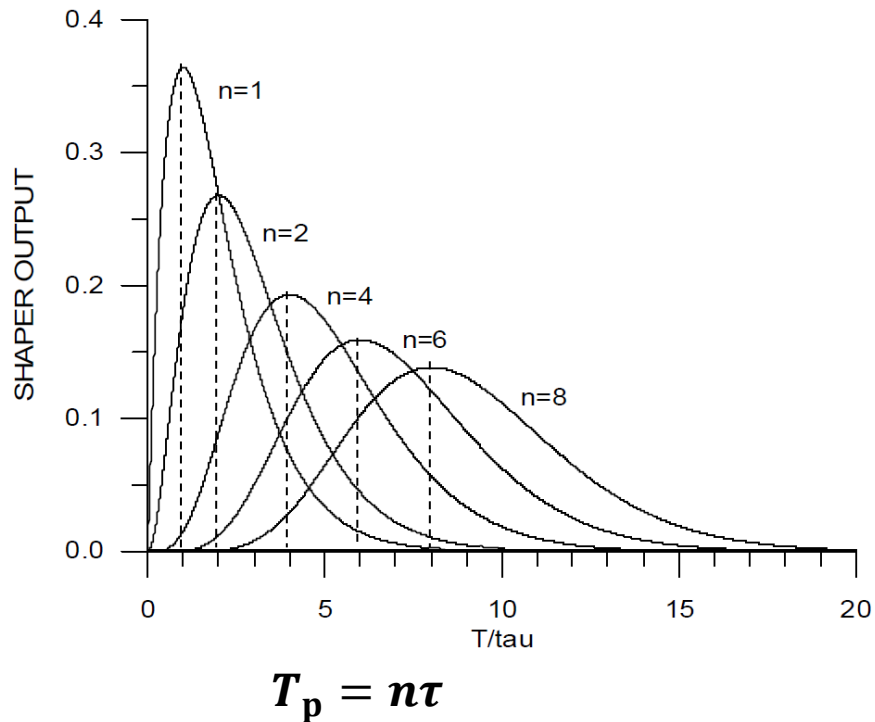
CR-RCⁿ pulse shaping combines a high-pass filter and n low-pass filter



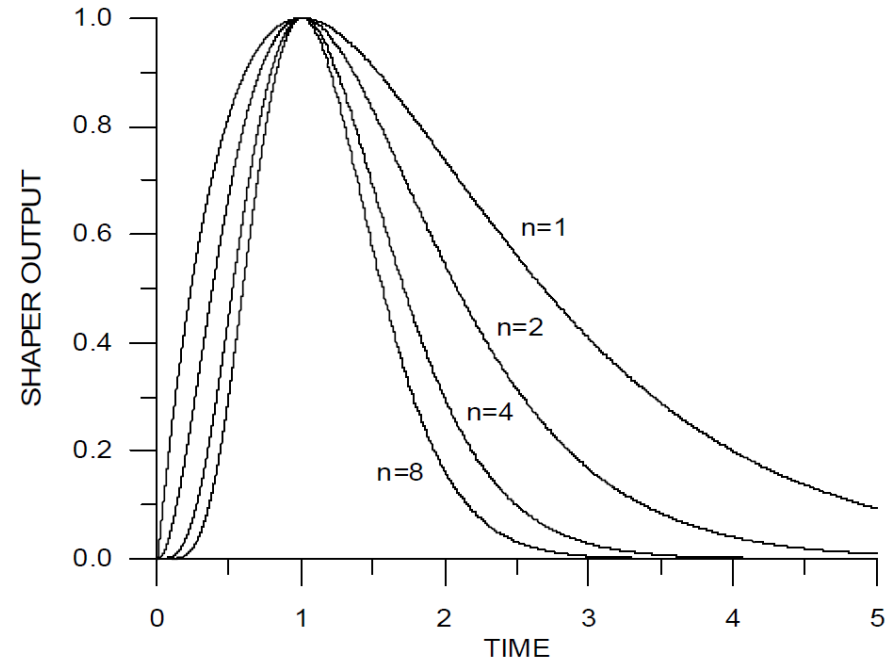
Pulses become gaussian-like after multiple integrators



- Integrating with the same time constant τ



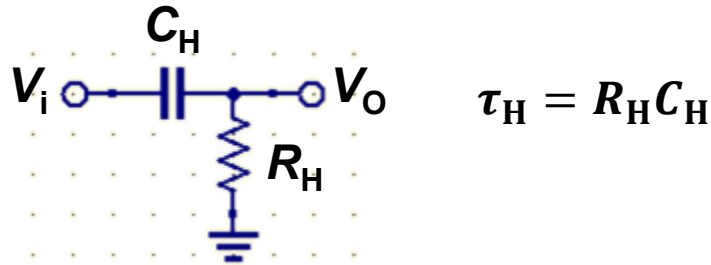
- Integrating with the reduced time constant $\tau_n = \frac{\tau_{n-1}}{n}$



CR is the high pass filter while RC is the low pass filter



- CR circuit:



$$V_o = V_i \frac{R_H}{\frac{1}{j\omega C_H} + R_H} = V_i \frac{1}{\frac{1}{j\omega R_H C_H} + 1}$$

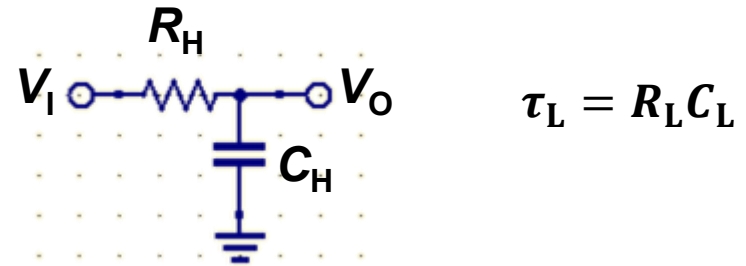
$$= V_i \frac{1}{\frac{1}{j\omega \tau_H} + 1}$$

For $\omega \tau_H \gg 1$, i. e.,

$$\frac{2\pi f}{f_H} \gg 1 \text{ or } f \gg f_H \text{ or } \tau_H \gg t$$

$$V_o \approx V_i \quad \bullet \text{ High-pass filter!}$$

- RC circuit:



$$V_o = V_i \frac{\frac{1}{j\omega C_L}}{\frac{1}{j\omega C_L} + R_L} = V_i \frac{1}{\frac{1}{j\omega R_L C_L} + 1}$$

$$= V_i \frac{1}{\frac{j\omega \tau_L}{1} + 1}$$

For $\omega \tau_L \ll 1$, i. e.,

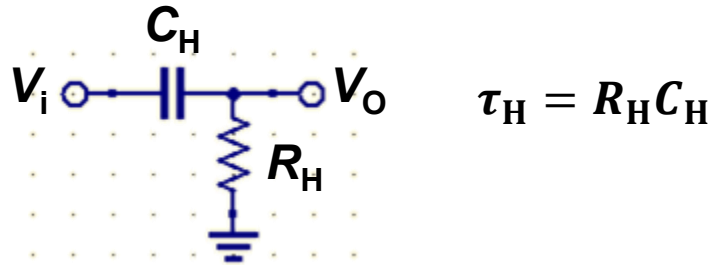
$$\frac{2\pi f}{f_L} \ll 1 \text{ or } f \ll f_L \text{ or } t \ll \tau_L$$

$$V_o \approx V_i \quad \bullet \text{ Low-pass filter!}$$

CR is the high pass filter while RC is the low pass filter



- CR circuit:



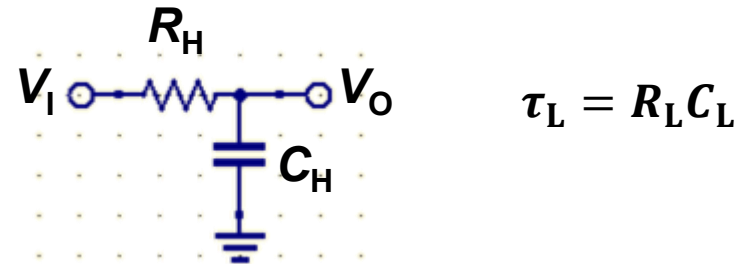
$$V_i - V_c - V_o = 0 \quad V_o = iR_H$$

$$V_c = \frac{1}{C_H} \int i dt = \frac{1}{C_H} \int \frac{V_o}{R_H} dt$$

$$\dot{V}_i = \frac{V_o}{\tau_H} + \dot{V}_o$$

$$V_o(t) = e^{-t/\tau_H} \int_0^t \dot{V}_i(t') e^{t'/\tau_H} dt'$$

- RC circuit:



$$V_i - V_R - V_o = 0 \quad i = C_L \dot{V}_o$$

$$V_R = iR_L = R_L C_L \dot{V}_o$$

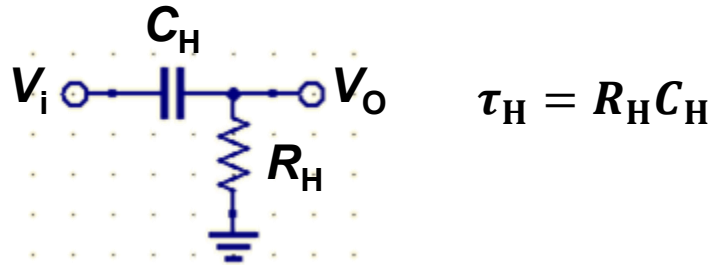
$$V_i = \tau_L \dot{V}_o + V_o$$

$$V_o(t) = e^{-t/\tau_L} \frac{1}{\tau_L} \int_0^t V_i(t') e^{t'/\tau_L} dt'$$

A step function becomes an exponential decay after the high-pass filter



- CR circuit:

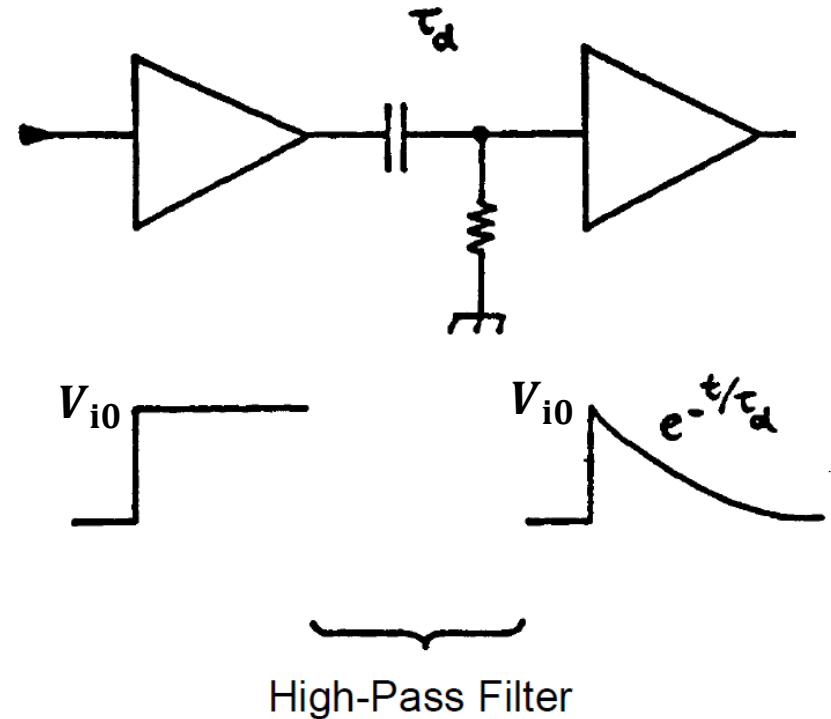


$$V_o(t) = e^{-t/\tau_H} \int_0^t \dot{V}_i(t') e^{t'/\tau_H} dt'$$

$$V_i(t) = \begin{cases} 0 & t < 0 \\ V_{i0} & t \geq 0 \end{cases}$$

$$\dot{V}_i(t) = V_{i0} \delta(t)$$

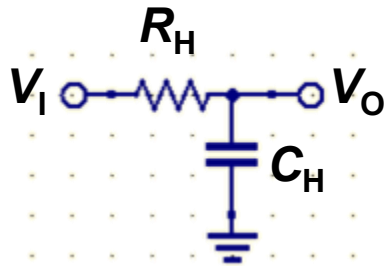
$$\begin{aligned} V_o(t) &= e^{-t/\tau_H} \int_0^t V_{i0} \delta(t') e^{t'/\tau_H} dt' \\ &= V_{i0} e^{-t/\tau_H} e^{t'/\tau_H} \Big|_{t'=0} = V_{i0} e^{-t/\tau_H} \end{aligned}$$



CR is the high pass filter while RC is the low pass filter



- RC circuit:



$$\tau_L = R_L C_L$$

$$V_o(t) = e^{-t/\tau_L} \frac{1}{\tau_L} \int_0^t V_i(t') e^{t'/\tau_L} dt'$$

$$V_i(t) = V_{i0} e^{-t/\tau_H}$$

For $t \ll \tau_L$ $e^{-t/\tau_L} \sim e^{t/\tau_L} \sim 1$

$$\begin{aligned} V_o(t) &\sim \frac{1}{\tau_L} \int_0^t V_i(t') dt' = \frac{1}{\tau_L} \int_0^t V_{i0} e^{-t'/\tau_H} dt' \\ &= \frac{\tau_H}{\tau_L} V_{i0} (1 - e^{-t/\tau_H}) \end{aligned}$$

$$V_o(t) = e^{-t/\tau_L} \frac{1}{\tau_L} \int_0^t V_{i0} e^{-t'/\tau_H} e^{t'/\tau_L} dt'$$

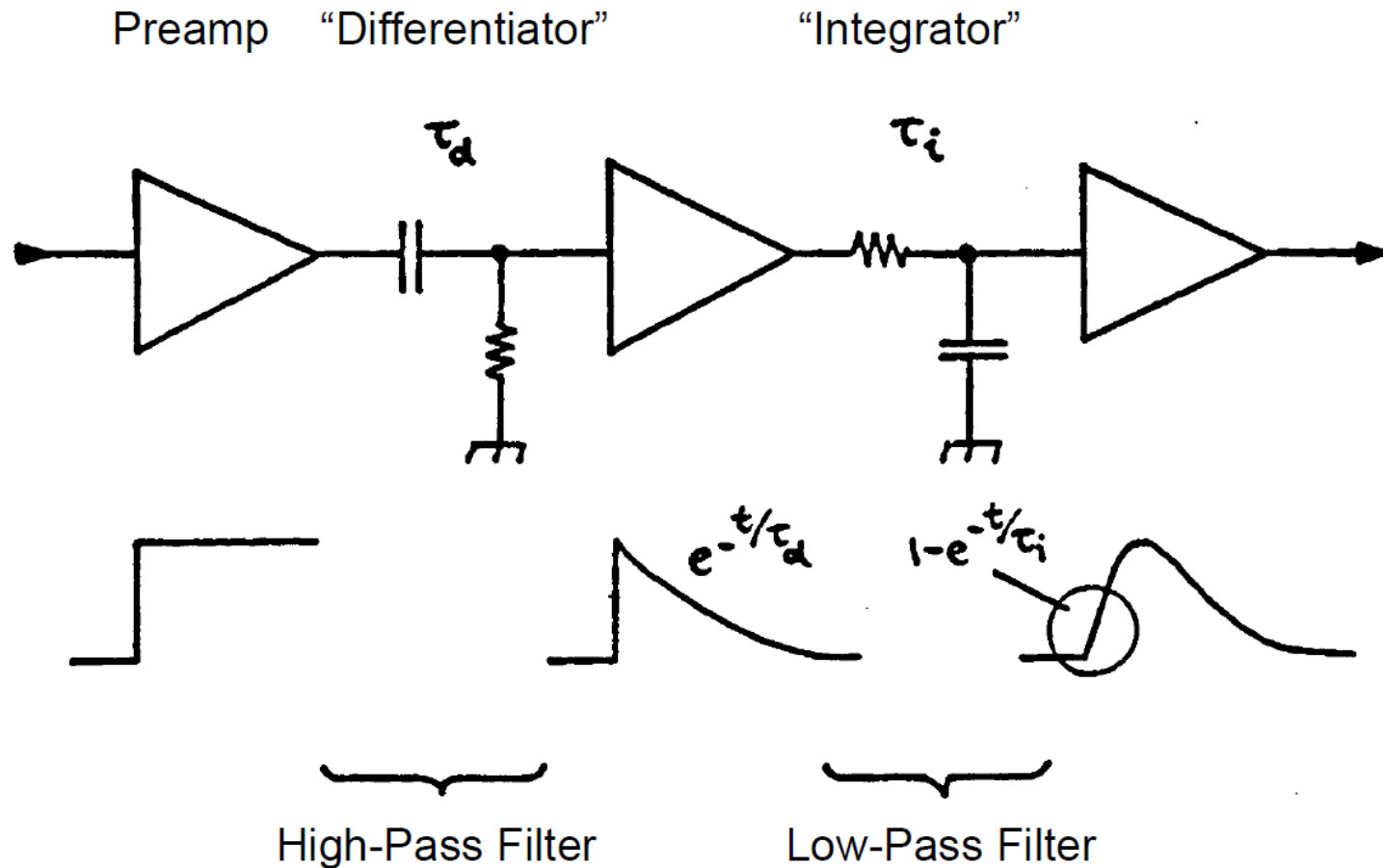
$$= e^{-t/\tau_L} \frac{1}{\tau_L} \int_0^t V_{i0} e^{t'/\tau_{\text{eff}}} dt'$$

$$= e^{-t/\tau_L} \frac{\tau_{\text{eff}}}{\tau_L} V_{i0} (e^{t/\tau_H} - 1)$$

$$\frac{1}{\tau_{\text{eff}}} = \frac{1}{\tau_L} - \frac{1}{\tau_H} < \frac{1}{\tau_L}$$



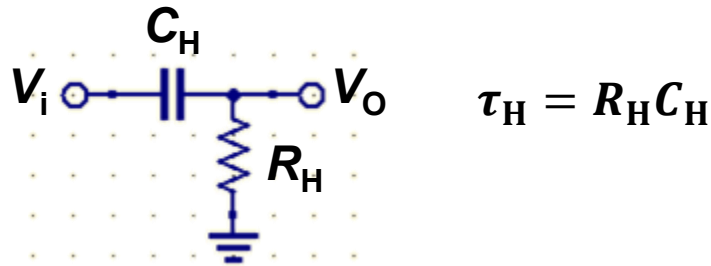
CR-RCⁿ pulse shaping combines a high-pass filter and n low-pass filter



Finite difference method is used to numerically applying the CR-RC pulse shaping



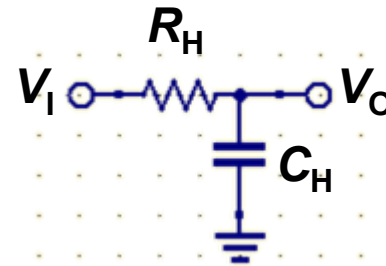
- CR circuit:



$$\tau_H = R_H C_H$$

$$\dot{V}_i = \frac{V_o}{\tau_H} + \dot{V}_o$$

- RC circuit:



$$\tau_L = R_L C_L$$

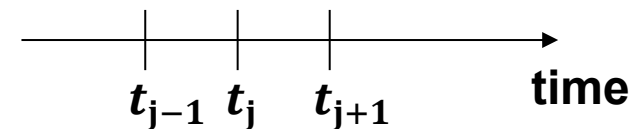
$$V_i = \tau_L \dot{V}_o + V_o$$

$$\begin{aligned} \frac{V_i(t_j) - V_i(t_{j-1})}{\Delta t} &\approx \frac{V_o(t_{j-1})}{\tau_H} + \frac{V_o(t_j) - V_o(t_{j-1})}{\Delta t} \\ &\approx \frac{V_o(t_j)}{\tau_H} + \frac{V_o(t_j) - V_o(t_{j-1})}{\Delta t} \end{aligned}$$

$$V_o(t_j) = \alpha_{CR} V_o(t_{j-1}) + \alpha_{CR} [V_i(t_j) - V_i(t_{j-1})]$$

$$\alpha_{CR} \equiv \frac{\frac{\tau_H}{\Delta t}}{\frac{\tau_H}{\Delta t} + 1} = \frac{\frac{R_H C_H}{\Delta t}}{\frac{R_H C_H}{\Delta t} + 1}$$

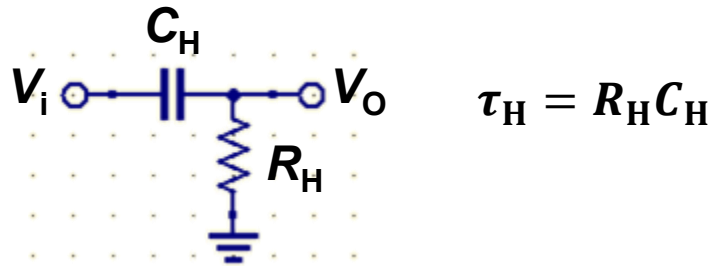
$$\begin{aligned} \frac{dV}{dt} &= \lim_{\Delta t \rightarrow 0} \frac{V(t_j) - V(t_{j-1})}{\Delta t} \\ &\approx \frac{V(t_j) - V(t_{j-1})}{\Delta t} \end{aligned}$$



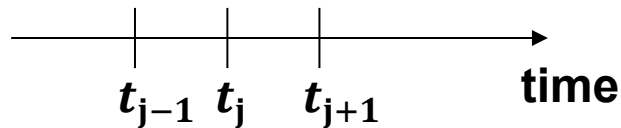
Finite difference method is used to numerically applying the CR-RC pulse shaping



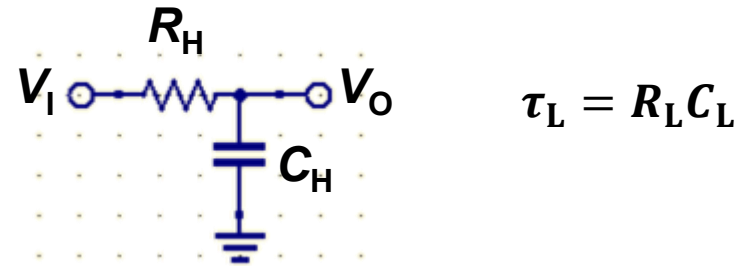
- CR circuit:



$$\dot{V}_i = \frac{V_o}{\tau_H} + \dot{V}_o$$



- RC circuit:



$$V_i = \tau_L \dot{V}_o + V_o$$

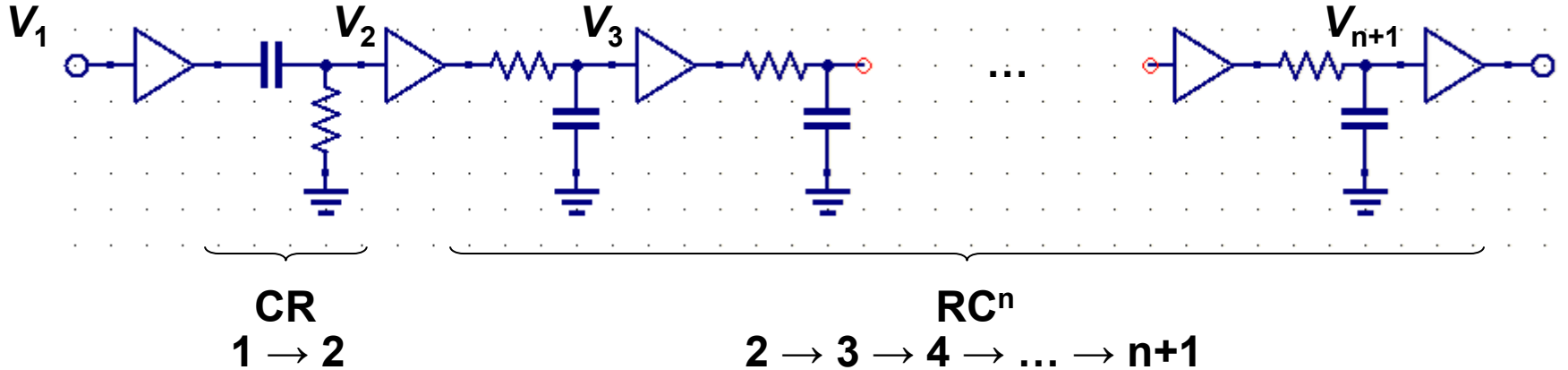
$$V_i(t_j) \approx \tau_L \frac{V_o(t_j) - V_o(t_{j-1})}{\Delta t} + V_o(t_{j-1})$$

$$\approx \tau_L \frac{V_o(t_j) - V_o(t_{j-1})}{\Delta t} + V_o(t_j)$$

$$V_o(t_j) = (1 - \alpha_{RC})V_o(t_{j-1}) + \alpha_{RC}V_i(t_j)$$

$$\alpha_{RC} \equiv \frac{1}{\frac{\tau_L}{\Delta t} + 1} = \frac{1}{\frac{R_L C_L}{\Delta t} + 1}$$

Iteration can be used to have more orders of CR-RCⁿ pulse shaping



$$V_2(t_j) = \alpha_{\text{CR}} V_2(t_{j-1}) + \alpha_{\text{CR}} [V_1(t_j) - V_1(t_{j-1})]$$

$$\alpha_{\text{CR}} \equiv \frac{\frac{\tau_H}{\Delta t}}{\frac{\tau_H}{\Delta t} + 1} = \frac{\frac{R_H C_H}{\Delta t}}{\frac{R_H C_H}{\Delta t} + 1}$$

$$V_3(t_j) = (1 - \alpha_{\text{RC},3}) V_3(t_{j-1}) + \alpha_{\text{RC},3} V_2(t_j)$$

$$\vdots$$

$$V_m(t_j) = (1 - \alpha_{\text{RC},m}) V_m(t_{j-1}) + \alpha_{\text{RC},m} V_{m-1}(t_j)$$

$$\vdots$$

$$V_{n+1}(t_j) = (1 - \alpha_{\text{RC},n+1}) V_{n+1}(t_{j-1}) + \alpha_{\text{RC},n+1} V_n(t_j)$$

$$\alpha_{\text{RC},m} \equiv \frac{1}{\frac{\tau_{L,m}}{\Delta t} + 1}$$

$$\tau_L = R_L C_L$$

$$\tau_{L,m} = \frac{\tau_{L,m-1}}{m}$$

Let's try using excel first



$$\alpha_{CR} \equiv 0.85$$

$$\alpha_{RC,2} \equiv 0.15 \quad \Delta t = 64 \text{ ns}$$

$$\alpha_{CR} \equiv \frac{\frac{\tau_H}{\Delta t}}{\frac{\tau_H}{\Delta t} + 1} = 0.85$$

$$\alpha_{RC,m} \equiv \frac{1}{\frac{\tau_{L,m}}{\Delta t} + 1} = 0.15$$

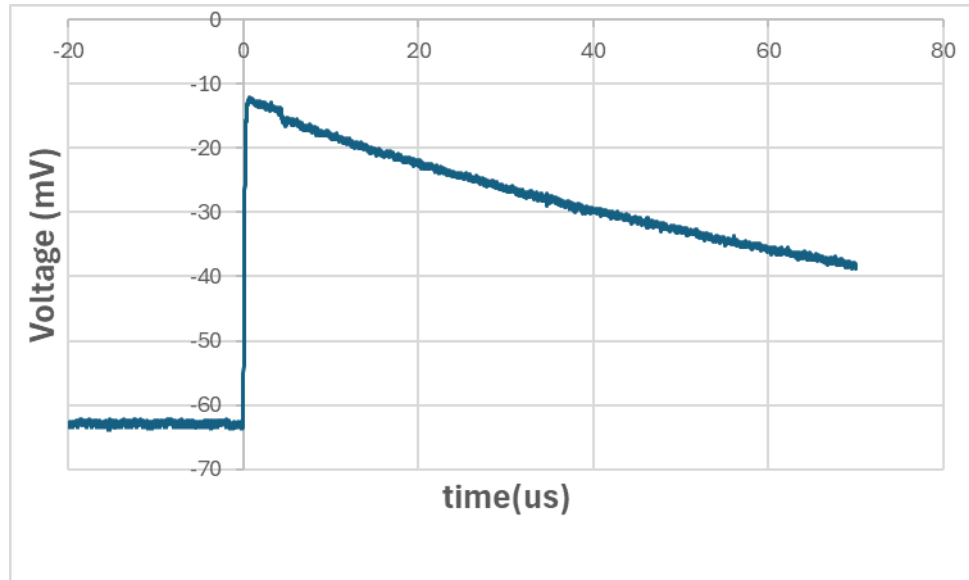
$$\tau_n = \frac{\tau_{n-1}}{n}$$

$$\tau_H = 362.7 \text{ ns}$$

$$\tau_{L,m} = 362.7/m \text{ ns}$$

$$\frac{\tau_{L,m}}{\Delta t} = \frac{0.85}{0.15}$$

$$\alpha_{RC,m} \equiv \frac{1}{\frac{0.85}{0.15 m} + 1}$$



$$V_2(t_j) = \alpha_{CR} V_2(t_{j-1}) + \alpha_{CR} [V_1(t_j) - V_1(t_{j-1})]$$

$$V_3(t_j) = (1 - \alpha_{RC,3}) V_3(t_{j-1}) + \alpha_{RC,3} V_2(t_j)$$

$$V_m(t_j) = (1 - \alpha_{RC,m}) V_m(t_{j-1}) + \alpha_{RC,m} V_{m-1}(t_j)$$

$$V_{n+1}(t_j) = (1 - \alpha_{RC,n+1}) V_{n+1}(t_{j-1}) + \alpha_{RC,n+1} V_n(t_j)$$

Let's try using excel first

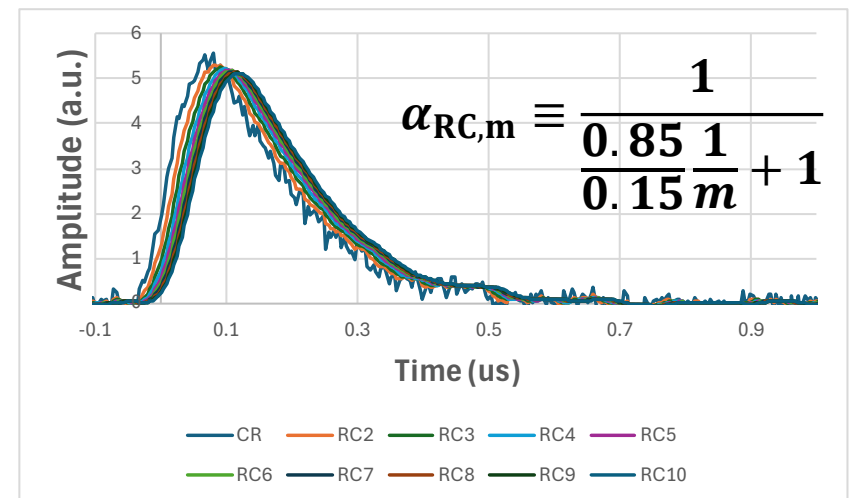
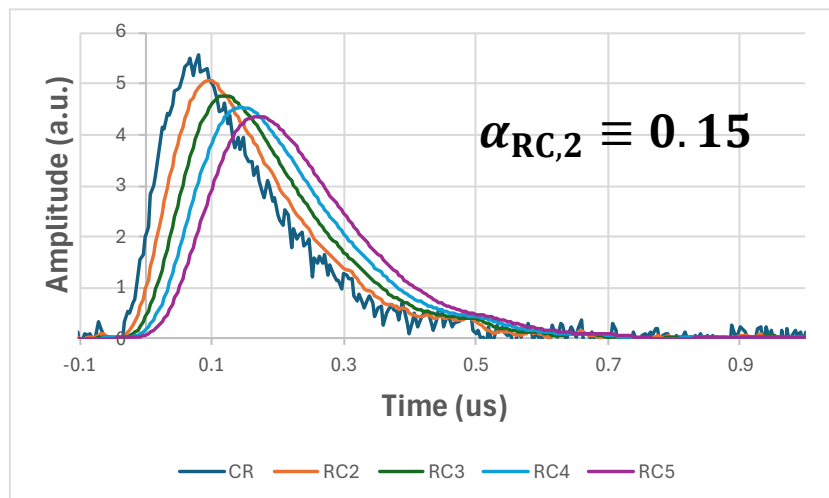
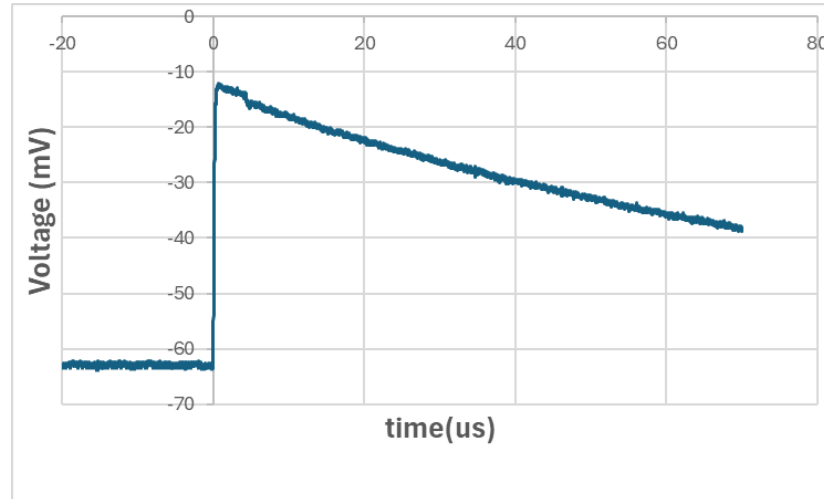


$$\Delta t = 64 \text{ ns}$$

$$\alpha_{\text{CR}} \equiv 0.85$$

$$\tau_{\text{H}} = 362.7 \text{ ns}$$

$$\tau_{\text{L},m} = 362.7/m \text{ ns}$$



When the pulse is not sharp enough, no significant gaussian feature is observed

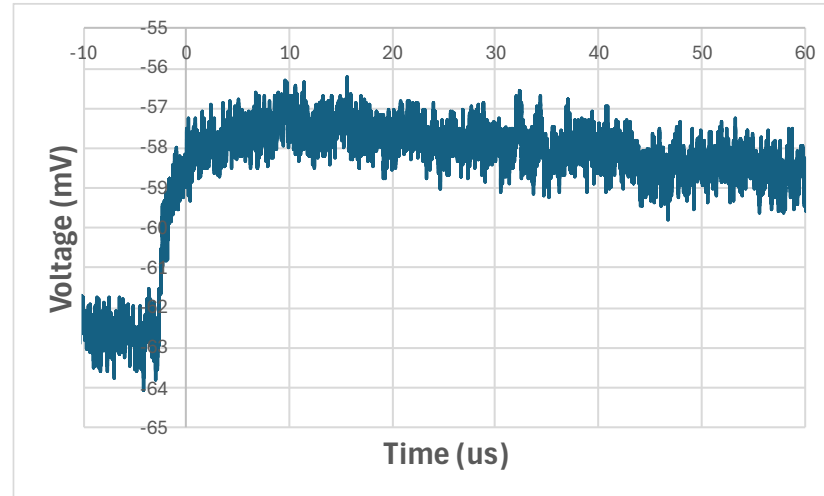


$$\Delta t = 64 \text{ ns}$$

$$\alpha_{\text{CR}} \equiv 0.85$$

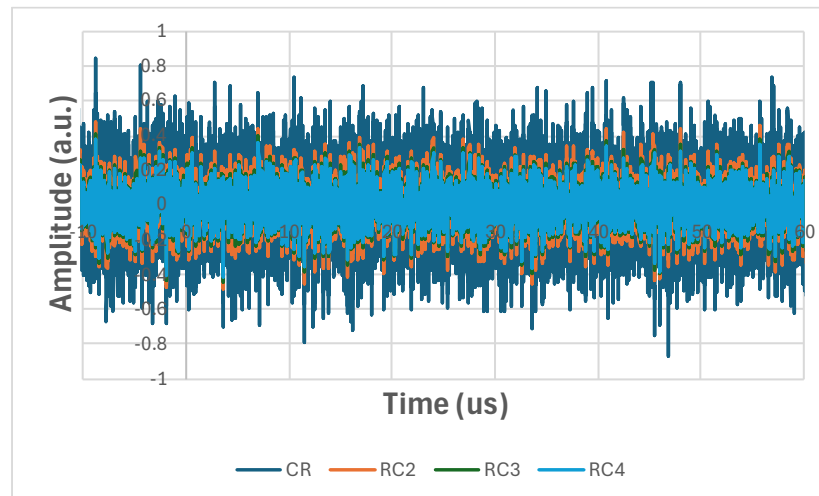
$$\tau_{\text{H}} = 362.7 \text{ ns}$$

$$\tau_{\text{L},m} = 362.7/m \text{ ns}$$



• What's the source?

$$\alpha_{\text{RC},m} \equiv \frac{1}{\frac{0.85}{0.15} \frac{1}{m} + 1}$$



Let's try using python



- You may install: IDLE / Visual Studio / Anaconda

```
1 import csv
2 acr = 0.85
3 arc = 0.15
4 brc = 1-arc
5 tau0vert = 0.85/0.15
6
7 filename = '20250610 (5)_0015'
8 fn2='{}.csv'.format(filename,j)
9 fn3='{}_f2.csv'.format(filename,j)
10 with open(fn2) as i:
11     with open(fn3,'x',encoding='utf-8', newline='') as o:
12         r = csv.reader(i)
13         w = csv.writer(o)
14         w.writerow(next(r))      # First three rows are header names
15         w.writerow(next(r))
16         w.writerow(next(r))
```

	A	B	C	D	E
1	time	voltage			
2	(us)	(mV)			
3					
4	-30.0079	-62.7857			
5	-30.0039	-62.9475			
6	-29.9999	-63.1062			
7	-29.9959	-62.7857			
8	-29.9919	-63.0543			
9	-29.9879	-63.1611			
10	-29.9839	-63.1062			
11	-29.9799	-62.5172			
12	-29.9759	-62.8407			

	A	B	C	D	E	F	G	H	I	J	K
1	time	voltage	CR	RC2	RC3	RC4	RC5				
2	(us)	(mV)									
3											
4	-30.0079	-62.7857	0	0	0	0	0				
5	-30.0039	-62.9475	-0.13749	-0.03587	-0.01242	-0.00514	-0.00241				
6	-29.9999	-63.1062	-0.25176	-0.09219	-0.04003	-0.01958	-0.01046				
7	-29.9959	-62.7857	0.058388	-0.05291	-0.04449	-0.02988	-0.01956				
8	-29.9919	-63.0543	0.17865	0.08571	0.05876	0.04183	0.03				

Let's try using python



```

17 first = next(r)
18 lastR = float(first[1])
19 last0 = 0 # lastR      # Later discovered that these starting
20 last1 = 0 # last0 * brc # values are all not very important
21 last2 = 0 # last1 * brc # so we can simply choose to start
22 last3 = 0 # last2 * brc # with all of them set to zero
23 last4 = 0 # last3 * brc # and things will still work fine
24 w.writerow( [first[0],lastR,last0,last1,last2,last3,last4])
25 for row in r:
26     last0 = acr * ( last0 + float(row[1]) - lastR )
27     last1 = (1-1/(tauOvert/2+1)) * last1 + 1/(tauOvert/2+1) * last0
28     last2 = (1-1/(tauOvert/3+1)) * last2 + 1/(tauOvert/3+1) * last1
29     last3 = (1-1/(tauOvert/4+1)) * last3 + 1/(tauOvert/4+1) * last2
30     last4 = (1-1/(tauOvert/5+1)) * last4 + 1/(tauOvert/5+1) * last3
31     lastR = float( row[1])
32     w.writerow( [ row[0],lastR,last0,last1,last2,last3,last4])

```

$$V_2(t_j) = \alpha_{CR} V_2(t_{j-1}) + \alpha_{CR} [V_1(t_j) - V_1(t_{j-1})]$$

$$V_m(t_j) = (1 - \alpha_{RC,m}) V_m(t_{j-1}) + \alpha_{RC,m} V_{m-1}(t_j)$$

$$\frac{\tau_{L,m}}{\Delta t} = \frac{0.85}{0.15}$$

$$\tau_n = \frac{\tau_{n-1}}{n}$$

$$\alpha_{RC,m} \equiv \frac{1}{\frac{0.85}{0.15} \frac{1}{m} + 1}$$

Course Outline



- **Formosa Integrated Research Spherical Tokamak (FIRST), the first Tokamak in Taiwan**
- **Physics used in Alpha-E**
 - **Microwave**
 - **ECRH**
 - **Spectrum of fusion products**
 - **Stopping power of the filter**
 - **Data analysis using Pulse Shaping**
- **Final report**

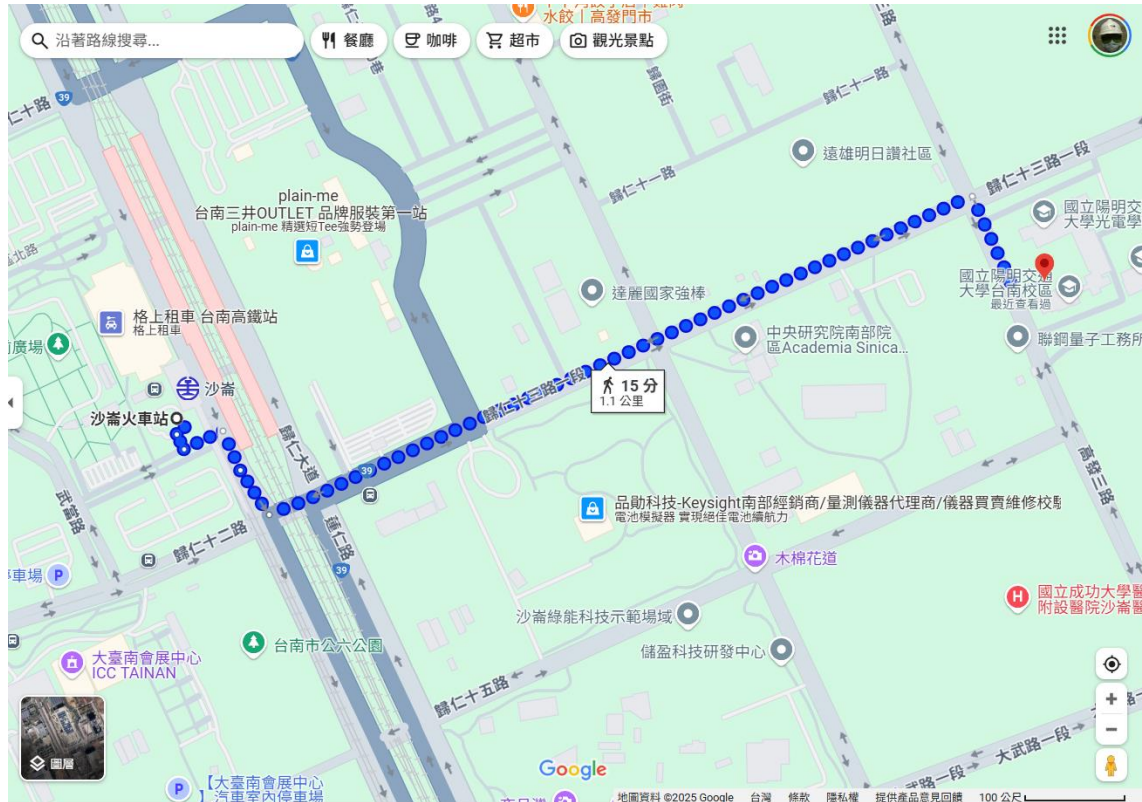
Final presentation



1. 實驗設備之原理。
2. 被分配的實驗項目的原理。
3. 數據分析原理說明。
4. 數據分析結果。
5. 討論。

若週四實驗課缺課，則報告為**0**分。

We will visit Alpha Ring's laboratory located in Shalun



- Alpha Ring @ Shalun
- **Meet at Tainan Station at 12:40.**
- Train 373, departing at 12:55.
- Walk to Alpha Ring arriving at 13:40.