Application of Plasma Phenomena



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Institute of Space and Plasma Sciences, National Cheng Kung University

2024 spring semester

Tuesday 9:10-12:00

Materials:

https://capst.ncku.edu.tw/PGS/index.php/teaching/

Online courses: https://nckucc.webex.com/nckucc/j.php?MTID=m4082f23c59af0571015416f6 e58dd803

2024/2/20 updated 1

Most of the material in space is plasma





Forty SpaceX's Starlink satellites were destroyed by a geomagnetic storm on 2022/2/4



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Geomagnetic storms occur when intense solar wind near Earth spawns shifting currents and plasmas in Earth's magnetosphere. This interaction can warm Earth's upper atmosphere and increase atmospheric density high enough above the planet to affect satellites in low orbits like SpaceX's new Starlink craft.

https://en.wikipedia.org/wiki/Geomagnetic_storm https://wonderfulengineering.com/watch-the-40-starlink-satellites-destroyed-by-a-geomagnetic-storm-burn-up-in-the-sky/

Plasma plays an important role on semiconductor manufacturing



 The process technology of Taiwan Semiconductor Manufacturing Company Limited (TSMC):



A semiconductor device is fabricated by many repetitive production process



Plasma is widely used in semiconductor fabrication



Deposition



EUV light source

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https://www.scorec.rpi.edu/research_plasmaetchmodeling.php http://Inf-wiki.eecs.umich.edu/wiki/Sputter_deposition

Plasma cleaning



R. S. Abhari, etc., J. Micro/Nanolithography, MEMS, and MOEMS, 11, 021114 (2012) 馗鼎奈米科技股份有限公司

https://www.ecplaza.net/products/plasma-cleaning_111807

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Nuclear fusion as an energy source is being developed



Magnetic confinement fusion (MCF)
 Inertial confinement fusion (ICF)





https://www.euro-fusion.org/2011/09/tokamak-principle-2/

Inertial confinement fusion: an introduction, Laboratory for Laser Energetics, University of Rochester

Nuclear fusion as an energy source is being developed



Magnetic confinement fusion (MCF)
 Inertial confinement fusion (ICF)



Significant breakthrough is achieved recently



Magnetic confinement fusion (MCF)
 Inertial confinement fusion (ICF)



 On 2024/2/(8), record-breaking 69.26 megajoules of sustained fusion energy in Joint European Torus (JET) facility in Oxford demonstrates powerplant potential and strengthens case for ITER.



 National Ignition Facility (NIF) demonstrated a gain grater than 1 for the first time on 2022/12/5. The yield of 3.15 MJ from the 2.05-MJ input laser energy, i.e., Q=1.5.

https://ccfe.ukaea.uk/resources/#gallery https://www.science.org/content/article/historic-explosion-long-sought-fusion-breakthrough

Plasma can be used as particle accelerators and thrusters



http://cuos.engin.umich.edu/researchgroups/hfs/research/laser-wakefield-acceleration/ https://zh.wikipedia.org/wiki/File:Electrostatic_ion_thruster-en.svg

V. Malka, *et al.*, Nature Physics **4**, 447 (2008)

Low temperature plasma is used in medical applications







Plasma medicine, by Alexander Fridman and Gary Friedman Biochem Biophys Res Commun. 2006 May 5; 343(2): 351–360.

Course Outline



- 1. What is Plasma?
- 2. Methods for Plasma Production
- 3. Planeterrella Artificial aurora demonstration
- 4. Plasma Diagnostics

Course Outline

- 5. Applications
 - 1. Material Processing
 - 2. Magnetron sputtering demonstration
 - 3. Plasma cleaning
 - 4. Light source and display systems
 - 5. Controlled thermonuclear fusion
 - 6. Magnetic mirror demonstration
 - 7. Plasma propulsion
 - 8. high energy particle accelerator
 - 9. biomedical application
 - 10. DBD discharge demonstration
- 6. Student presentations

Course Outline - Demonstration

- a. Planeterrella
- b. Magnetron sputtering
- c. Dielectric barrier discharge (DBD)
- d. Magnetic mirror
- e. Tesla coil



Planeterrella



Magnetron sputtering



DBD plasma







Tesla Coil



- Quizzes 50 % (2-min Q&A at the beginning of each class)
- Presentations 50 % (10-min presentation on any

plasma applications or phenomena)



Reference for the section "What is Plasma?"

- Introduction to plasma theory, by Dwight R. Nicholson
- Introduction to plasma physics and controlled fusion, by Francis F. Chen
- Principles of plasma physics for engineers and scientists, by Umran S.
 Inan and Marek Golkowski
- The physics of plasma, by T. J. M. Boyd and J. J. Sanderson
- Principles of plasma physics, by Krall and Trivelpiece
- NRL Plasma Formulary, Naval Research Laboratory, 2013 by J. D. Huba

Course Outline



1. What is Plasma?

- 2. Methods for Plasma Production
- 3. Planeterrella Artificial aurora demonstration
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Charged particles are accelerated due to Lorentz force under electromagnetic fields

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- Lorentz force: $\overrightarrow{F} = q \, \overrightarrow{E} + q \, \overrightarrow{v} \times \overrightarrow{B}$
 - -
- Force under electric fields







Charged particles gyro around magnetic field lines





http://www.ipp.cas.cz/vedecka_struktura_ufp/tokamak/tokamak_compass/diagnostics/ mikrovInne-diagnostiky/ece-ebw-radiometr.html

http://www-ssg.sr.unh.edu/tof/Smart/Students/lees/periods.html

Plasma is the 4th state of matter





http://tetronics.com/our-technology/what-is-plasmado

Plasma is everywhere







https://lasers.llnl.gov/science/understanding-the-universe/plasma-physics http://lnf-wiki.eecs.umich.edu/wiki/Sputter_deposition https://simple.wikipedia.org/wiki/Fluorescent lamp

In plasma, there are ions, electrons, and neutral gas



A plasma is a gas in which an important fraction of the atoms is ionized so that the electrons and ions are separated freely



http://ocw.mit.edu/courses/nuclear-engineering/22-611j-introduction-to-plasma-physics-i-fall-2003/lecture-notes/

A plasma can be created when the ionization rate is higher than the recombination rate



J. D. Huba \NRL Plasma Formulary", Naval Research Laboratory, 20134

There are several Important plasma parameters that need to be considered



• Debye length

$$\lambda_{\rm D} \equiv \left(\frac{KT_{\rm e}}{4\pi ne^2}\right)^{1/2}$$

Plasma parameter

$$\Lambda \equiv n \frac{4\pi}{3} \lambda_{\rm D}^{3}$$

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Plasma frequency

$$\omega_{\rm pe} \equiv \left(\frac{4\pi n_{\rm e}e^2}{m_{\rm e}}\right)^{1/2}$$

• Collision time $au_{\rm e} \equiv \frac{3\sqrt{m_{\rm e}}(KT_{\rm e})^{3/2}}{4\sqrt{2\pi}ne^4 \ln\Lambda}$

• Hall parameter $\chi \equiv \omega_{ce} \tau_{e}$, where $\omega_{ce} \equiv \frac{eB}{m_{e}c}$ is the electron gyrofrequency

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• Plasma beta $\beta \equiv \frac{P}{P_{\rm B}}$, where $P_{\rm B} \equiv \frac{B^2}{8\pi}$ is the magnetic pressure

A test ion in the plasma gathers a shielding cloud that tends to cancel its own charge



Francis F. Chen, \Introduction to plasma physics and controlled fusion²₂₆

Debye shielding is a phenomenon such that the potential due to a test charge in a plasma falls off much faster than in vacuum



• Vacuum potential:

$$\phi = \frac{\phi_0}{r}$$

Poisson's equation:

$$\nabla^2 \phi = 4\pi e(n_{\rm e} - n_{\rm i}) - 4\pi q_{\rm T} \delta(\vec{r})$$

Density profile:

$$n_{\rm e} = n_0 \exp\left(\frac{e\phi}{KT_{\rm e}}\right)$$
 , $n_{\rm i} = n_0$

For $\vec{r} \neq 0$ and assuming $\frac{e\phi}{T_e} \ll 1$

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\Phi}{dr} \right) = \frac{4\pi n_0 e^2}{KT_e} \phi$$
$$\phi = \frac{\phi_0}{r} \exp\left(-\frac{r}{\lambda_D}\right) \quad \lambda_D \approx \left(\frac{KT_e}{4\pi n e^2}\right)^{1/2}$$

Plasma parameter Λ is the number of particles in a sphere with radius of λ_D





• Plasma parameter:

$$\Lambda \equiv n \frac{4\pi}{3} \lambda_{\rm D}^3$$

• Criterion for an ionized gas to be plasma:

 $\lambda_{\rm D} \ll L$

• Requirement of "collective behavior":

 $\Lambda \gg 1$

Electron plasma frequency is the characteristic frequency such that electrons oscillate around their equilibrium positions



• Assumption:

$$\overrightarrow{\nabla} \equiv \widehat{x} \frac{\partial}{\partial x}, \ \overrightarrow{E} = \widehat{x}E, \ \overrightarrow{v}_{e} = \widehat{x}v,$$
,
 $\overrightarrow{\nabla} \times \overrightarrow{E} = 0, \ \overrightarrow{E} = -\overrightarrow{\nabla}\phi$

• Continuity and momentum equation for electron:

$$\frac{\partial n_{\rm e}}{\partial t} + \vec{\nabla} \cdot (n_{\rm e} \vec{v}_{\rm e}) = 0$$
$$m_{\rm e} n_{\rm e} \left[\frac{\partial \vec{v}_{\rm e}}{\partial t} + (\vec{v}_{\rm e} \cdot \vec{\nabla}) \vec{v}_{\rm e} \right] = -e n_{\rm e} \vec{E}$$

- Gauss' law:
 - $\frac{\partial E}{\partial x} = 4\pi e(n_{\rm i} n_{\rm e})$

Electron plasma frequency is obtained by linearizing the hydrodynamic equations

The oscillation is assumed to be small:

$$n_{\mathrm{e}}=n_{0}+n_{1},\,\overline{E}=\overline{E}_{0}+\overline{E}_{1}\,,\,v_{\mathrm{e}}=v_{0}+v_{1}$$
 where

$$\frac{\partial n_0}{\partial x} = v_0 = \vec{E}_0 = 0$$
$$\frac{\partial n_0}{\partial t} = \frac{\partial v_0}{\partial t} = \frac{\partial \vec{E}_0}{\partial t} = 0$$

Linearization:

$$m_{e}\frac{\partial v_{1}}{\partial t} = -eE_{1}$$
$$\frac{\partial n_{1}}{\partial t} + n_{0}\frac{\partial v_{1}}{\partial x} = 0$$
$$\frac{\partial E_{1}}{\partial x} = -4\pi en_{1}$$

Plane wave solution:

$$\eta_1 = \widehat{\eta} \exp[i(kx - \omega t)]$$

$$\eta_1 = v_1, n_1, E_1$$

• Substitute into the previous equations:

$$-im_{e}\omega v_{1} = -eE_{1}$$
$$-i\omega n_{1} = -n_{0}ikv_{1}$$
$$ikE_{1} = -4\pi en_{1}$$

• Electron plasma frequency is obtained by eliminating n₁ and E₁: $\omega_{pe} \equiv \omega = \left(\frac{4\pi n_e e^2}{m_o}\right)^{1/2}$

Electromagnetic wave can not propagate in a plasma with density higher than critical density



• Linearization:

$$\overrightarrow{\nabla} \times \overrightarrow{E}_{1} = -\frac{1}{c} \partial_{t} \overrightarrow{B}_{1}$$
$$\overrightarrow{\nabla} \times \overrightarrow{B}_{1} = \frac{4\pi e}{c} n_{0} \overrightarrow{v}_{e} + \frac{1}{c} \partial_{t} \overrightarrow{E}_{1}$$
$$m_{e} n_{0} \partial_{t} \overrightarrow{v}_{1} = -e n_{0} \overrightarrow{E}_{1}$$

• Eliminate
$$\partial_t \vec{v}_1$$
 and \vec{B}_1

$$-c \overrightarrow{\nabla} \times \left(\overrightarrow{\nabla} \times \overrightarrow{E}_{1} \right) = \frac{1}{c} \frac{4\pi n_{0} e^{2}}{m_{e}} \overrightarrow{E}_{1} + \frac{1}{c} \frac{\partial^{2} \overrightarrow{E}_{1}}{\partial t^{2}}$$

Take
$$\overrightarrow{\nabla} \times (\overrightarrow{\nabla} \times) \rightarrow k^2$$
 and $\partial_t \rightarrow \omega$
 $(\omega^2 - k^2 c^2 - \omega_e^2) \overrightarrow{E}_1 = 0$
 $\omega^2 - k^2 c^2 - \omega_e^2 = 0$

• Wave number k becomes imaginary when $\omega < \omega_e$.

The cutoff of the electromagnetic wave is important in laser fusion and in the interaction of radio waves with the ionosphere



SpaceX moves their S-band transmitter to the top of their rocket to avoid communication blackout





Charged particles collide with each other through collisions





$$\mathbf{m}\boldsymbol{v}_{\perp} = \int_{-\infty}^{\infty} dt \, \mathbf{F}_{\perp}(t)$$

Coulomb force:

$$m\,\frac{\ddot{r}}{r}=\frac{qq_0}{r^2}\hat{r}$$

$$F_{\perp} = \frac{qq_0}{p^2} \sin^3\theta$$

• Relation between θ and t is

$$x = -r\cos\theta = -\frac{p\cos\theta}{\sin\theta} = v_0 t$$

• Therefore,

$$v_{\perp} = \frac{qq_0}{mv_0p} \int_0^{\pi} d\theta \sin\theta = \frac{2qq_0}{mv_0p} \equiv \frac{v_0p_0}{p}$$

where
$$p_0 \equiv \frac{2qq_0}{m{v_0}^2}$$

• Note that this is valid only when $v_{\perp} << v_0$, i.e., $p >> p_0$.

Cumulative effect of many small angle collisions is more important than large angle collisions

• Consider a variable Δx that is the sum of many small random variables Δx_i , i=1,2,3,...,N,

$$\Delta x = \Delta x_1 + \Delta x_2 + \Delta x_3 + \dots + \Delta x_N = \sum_{i=1}^{N} \Delta x_i$$

< $\Delta x_i \ge < \Delta x_i \Delta x_i \ge \ldots = 0$

• Suppose $<\Delta x_i>=<\Delta x_i\Delta x_j>_{i\neq j}=0$

$$\langle (\Delta x)^2 \rangle = \left| \left(\sum_{i=1}^N \Delta x_i \right)^2 \right| = \sum_{i=1}^N \langle (\Delta x_i)^2 \rangle = N \langle (\Delta x_i)^2 \rangle$$

• For one collision:

$$\left\langle \boldsymbol{v}_{\perp}^{2} \right\rangle = \left\langle (\boldsymbol{\Delta}\boldsymbol{v}_{\mathrm{x}})^{2} \right\rangle + \left\langle \left(\boldsymbol{\Delta}\boldsymbol{v}_{\mathrm{y}} \right)^{2} \right\rangle = \frac{\boldsymbol{v}_{0}^{2} \boldsymbol{p}_{0}^{2}}{\boldsymbol{p}^{2}} \qquad \left\langle (\boldsymbol{\Delta}\boldsymbol{v}_{\mathrm{x}})^{2} \right\rangle = \left\langle \left(\boldsymbol{\Delta}\boldsymbol{v}_{\mathrm{y}} \right)^{2} \right\rangle = \frac{1}{2} \frac{\boldsymbol{v}_{0}^{2} \boldsymbol{p}_{0}^{2}}{\boldsymbol{p}^{2}}$$

• The total velocity in \hat{x}

$$\left\langle \left(\Delta v_{\mathrm{x}}^{\mathrm{tot}} \right)^{2} \right\rangle = N \left\langle (\Delta v_{\mathrm{x}})^{2} \right\rangle = \frac{N}{2} \frac{v_{0}^{2} p_{0}^{2}}{p^{2}}$$

The collision frequency can be obtained by integrating all the possible impact parameter



• Number of collisions in a time interval:

$$dN = n_0 2\pi p \, dp \, v_0 \, dt$$

.e., $\frac{dN}{dt} = 2\pi p \, dp \, n_0 v_0$

• Therefore

$$\frac{d}{dt}\left\langle \left(\Delta v_{x}^{\text{tot}}\right)^{2}\right\rangle = \frac{1}{2}\frac{v_{0}^{2}p_{0}^{2}}{p^{2}}\frac{dN}{dt}$$
$$= \pi n_{0}v_{0}^{3}p_{0}^{2}\frac{dp}{p}$$

$$\frac{d}{dt} \left\langle \left(\Delta_{\perp}^{\text{tot}} \right)^2 \right\rangle = 2 \frac{d}{dt} \left\langle \left(\Delta v_x^{\text{tot}} \right)^2 \right\rangle$$
$$= 2 \pi n_0 v_0^3 p_0^2 \int_{p_{\min}}^{p_{\max}} \frac{dp}{p}$$
$$= 2 \pi n_0 v_0^3 p_0^2 \ln \left(\frac{p_{\max}}{p_{\min}} \right)$$
$$\approx 2 \pi n_0 v_0^3 p_0^2 \ln \left(\frac{\lambda_D}{|p_0|} \right)$$
$$\approx 2 \pi n_0 v_0^3 p_0^2 \ln \Lambda$$

• Note that $\lambda_{\rm D} \approx \left(\frac{KT_{\rm e}}{4\pi n_0 e^2}\right)^{1/2}$ $\frac{\lambda_{\rm D}}{|p_0|} \approx \frac{\lambda_{\rm D} m_{\rm e} v_{\rm e}^2}{2e^2} \approx \frac{\lambda_{\rm D} KT_{\rm e}}{e^2} \approx 4\pi n_0 \lambda_{\rm D}^3$ $\approx \Lambda$

Comparison between the mean free path and the system size L determines the regime of the plasma

• A reasonable definition for the scattering time due to small angle collisions is the time it takes $\langle (\Delta v_{\perp}^{tot})^2 \rangle$ to equal v_0^2 . The collision frequency v_c due to small-angle collisions:

$$\frac{d}{dt} \left\langle \left(\Delta_{\perp}^{\text{tot}} \right)^2 \right\rangle \approx 2\pi n_0 v_0^3 p_0^2 \ln \Lambda \approx v_0^2 v_c, \quad p_0 \equiv \frac{2qq_0}{m_e v_0^2} \Rightarrow v_c = \frac{8\pi n_0 e^4 \ln \Lambda}{m_e^2 v_0^3}$$

• With more careful derivation, the collisional time is obtained:

$$\tau_{\rm e}^{-1} = \nu_{\rm c} = \frac{4\sqrt{2\pi}ne^4\ln\Lambda}{3\sqrt{m_{\rm e}}(KT_{\rm e})^{3/2}}$$

Mean free path:

 $l_{\rm mfp} = v_{\rm e} \tau_{\rm e}$

 $\begin{cases} l_{mfp} < L & Fluid Theory \\ l_{mfp} > L & Kinetic Theory \end{cases}$

Thermal conduction perpendicular to the magnetic field can be suppressed when the plasma is magnetized



Plasma is magnetized when

$$\frac{R_{\rm L}}{l_{\rm mfp}} = \frac{v_{\rm e}}{\omega_{\rm ce}} \frac{1}{v_{\rm e}\tau_{\rm e}} < 1$$

i.e., the hall parameter

 $\chi \equiv \omega_{\rm ce} \tau_e > 1$

$$m_{\rm e}\frac{d\,\vec{v}}{dt}=-\frac{e}{c}\,\vec{v}\times\vec{B}$$

• Assuming $\overrightarrow{B} = B\widehat{z}$ and the electron oscillates in x-y plane

$$m_{\rm e}v_{\rm x} = -\frac{e}{c}Bv_{\rm y} \qquad m_{\rm e}v_{\rm z} = 0$$
$$m_{\rm e}v_{\rm y} = \frac{e}{c}Bv_{\rm x}$$

$$\ddot{v}_{x} = -\frac{eB}{m_{e}c}\dot{v}_{y} = -\left(\frac{eB}{m_{e}c}\right)^{2}v_{x}$$
$$\ddot{v}_{y} = -\frac{eB}{m_{e}c}\dot{v}_{x} = -\left(\frac{eB}{m_{e}c}\right)^{2}v_{y}$$

• Therefore

$$\omega_{\rm ce} = \frac{eB}{m_{\rm e}c}$$

Plasma β is the ratio between hydro pressure and magnetic pressure



• Momentum equation in Magnetohydrodynamics (MHD) approach:

$$\rho \frac{d \vec{v}}{dt} + \rho \left(\vec{v} \cdot \vec{\nabla} \right) \vec{v} = - \vec{\nabla} p + \frac{1}{c} \vec{j} \times \vec{B}$$
$$\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j}$$

$$\vec{j} \times \vec{B} = \frac{c}{4\pi} \left(\vec{\nabla} \times \vec{B} \right) \times \vec{B} = \frac{c}{4\pi} \left[\left(\vec{B} \cdot \vec{\nabla} \right) \vec{B} - \frac{1}{2} \vec{\nabla} B^2 \right]$$
$$\rho \frac{d \vec{v}}{dt} + \rho \left(\vec{v} \cdot \vec{\nabla} \right) \vec{v} = - \vec{\nabla} \left(p + \frac{B^2}{8\pi} \right) + \frac{1}{4\pi} \left(\vec{B} \cdot \vec{\nabla} \right) \vec{B}$$

 $\frac{B^2}{8\pi}$

- Magnetic pressure:
- Magnetic tension: $\frac{1}{4\pi} (\vec{B} \cdot \vec{\nabla}) \vec{B}$

 $\beta \equiv \frac{p}{B^2/8\pi}$

There are several Important plasma parameters that need to be considered



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Plasma parameter

$$\Lambda \equiv n \frac{4\pi}{3} \lambda_{\rm D}^{3}$$

Plasma frequency

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• Hall parameter $\chi \equiv \omega_{ce} \tau_{e}$, where $\omega_{ce} \equiv \frac{eB}{m_{e}c}$ is the electron gyrofrequency

2

• Plasma beta $\beta \equiv \frac{P}{P_{\rm B}}$, where $P_{\rm B} \equiv \frac{B^2}{8\pi}$ is the magnetic pressure